



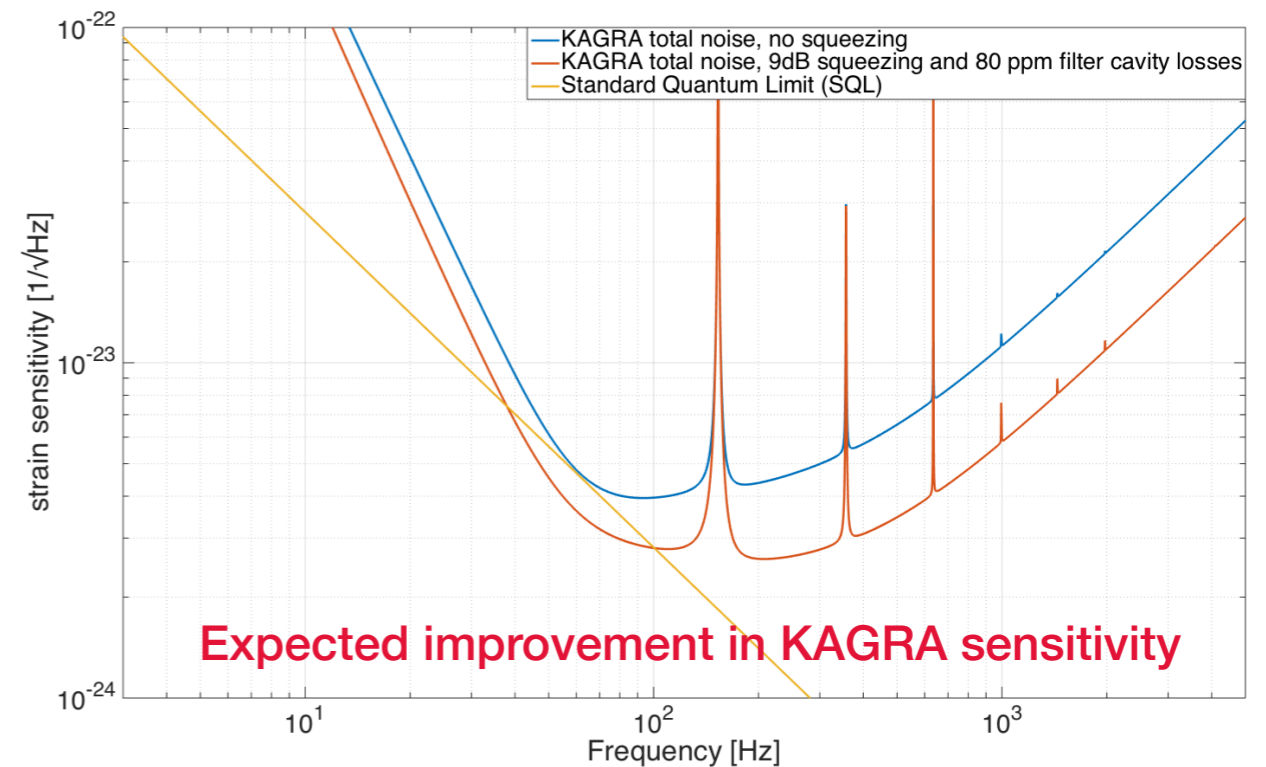
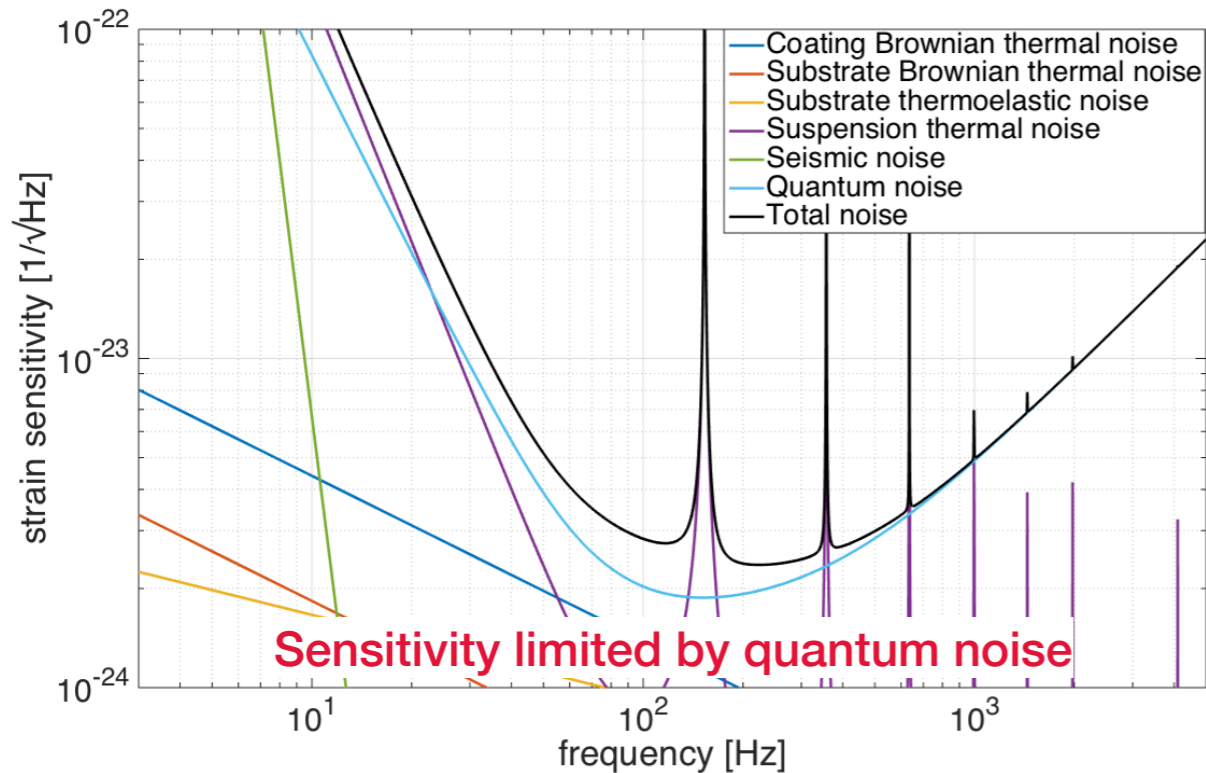
Status of the frequency dependent squeezing experiment at TAMA

K. Arai, Y. Aso, M. Barsuglia, E. Capocasa, M. Eisenmann, R. Flaminio, Y. Guo, M. Leonardi, M. Marchiò, L. Pinard, P. Prat, R. Schnabel, E. Schreiber, K. Somiya, M. Tacca, R. Takahashi, D. Tatsumi, A. Tomura, M. Vardaro, Y. Zhao

KAGRA F2F meeting, Toyama University, 24-26 August 2018

Motivation

- Broadband quantum noise reduction



PHYSICAL REVIEW D 93, 082004 (2016)

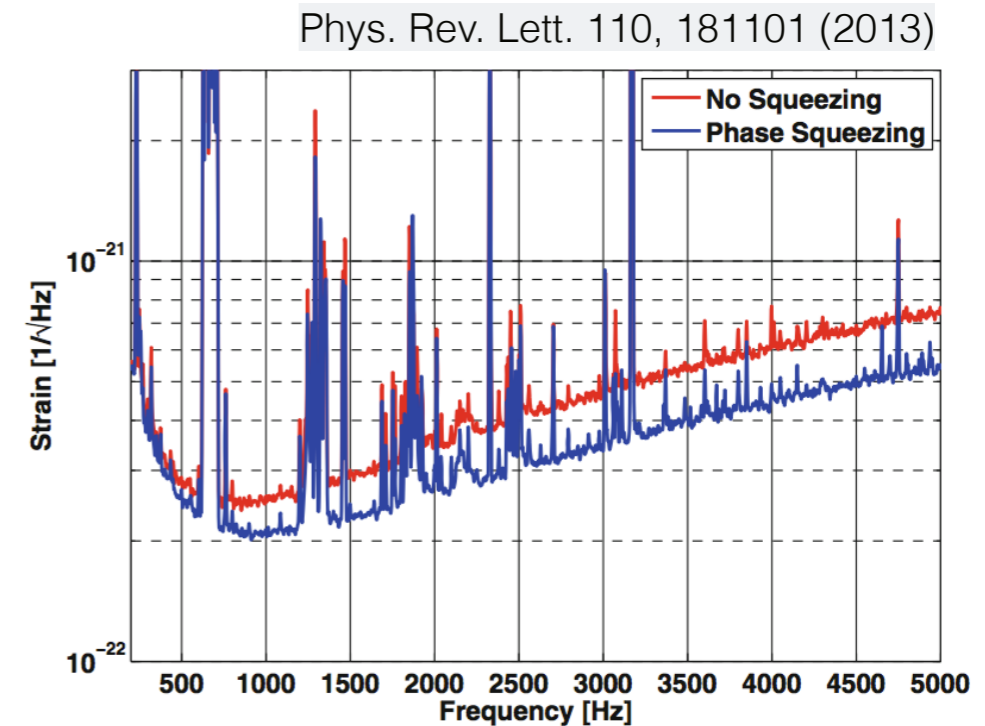
Estimation of losses in a 300 m filter cavity and quantum noise reduction in the KAGRA gravitational-wave detector

300 m filter cavity planned for ALIGO and AdVIRGO in O4

Frequency (in)dependent squeezing effect

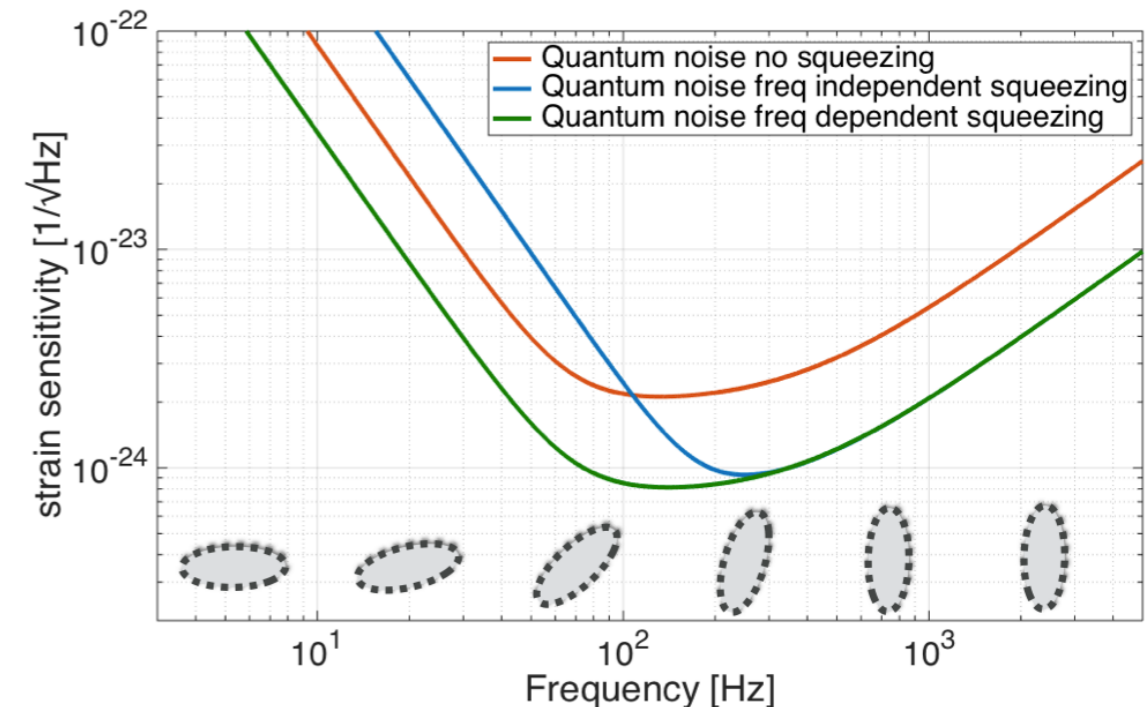
Frequency independent squeezing:

- Already successfully tested in GEO and LIGO
- Installed in LIGO and Virgo for O3
- Improvements at high frequency only



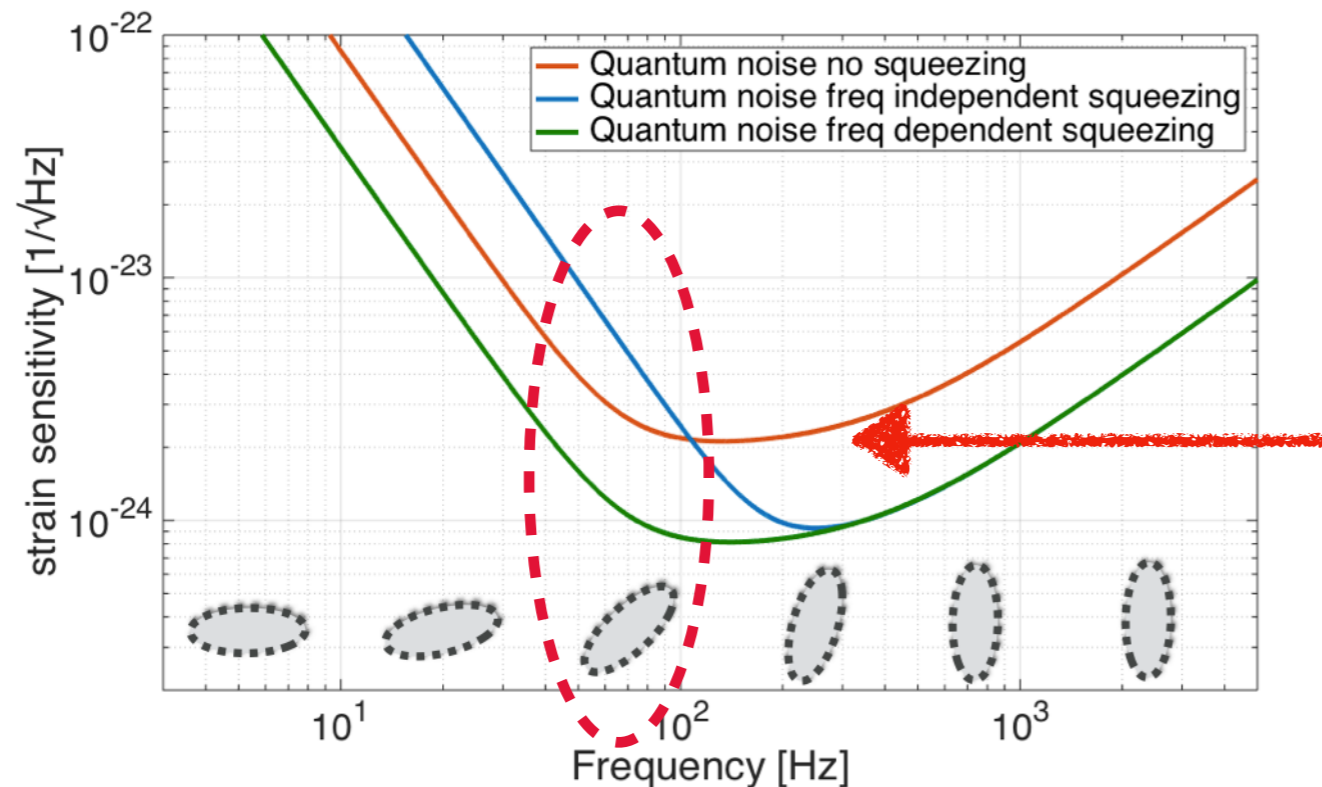
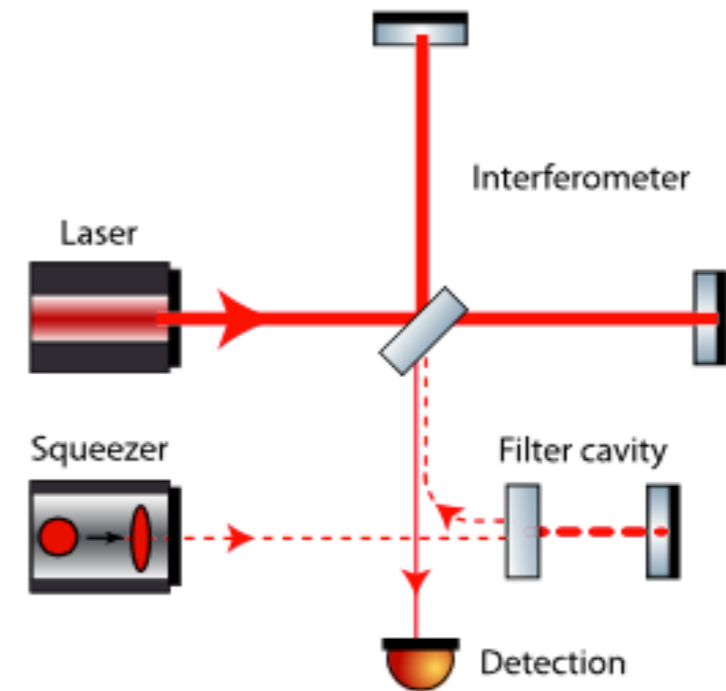
Frequency dependent squeezing:

- Counteract IFO squeezing rotation
- Planned for O4 in LIGO and Virgo
- Broadband quantum noise reduction



How to produce frequency dependent squeezing?

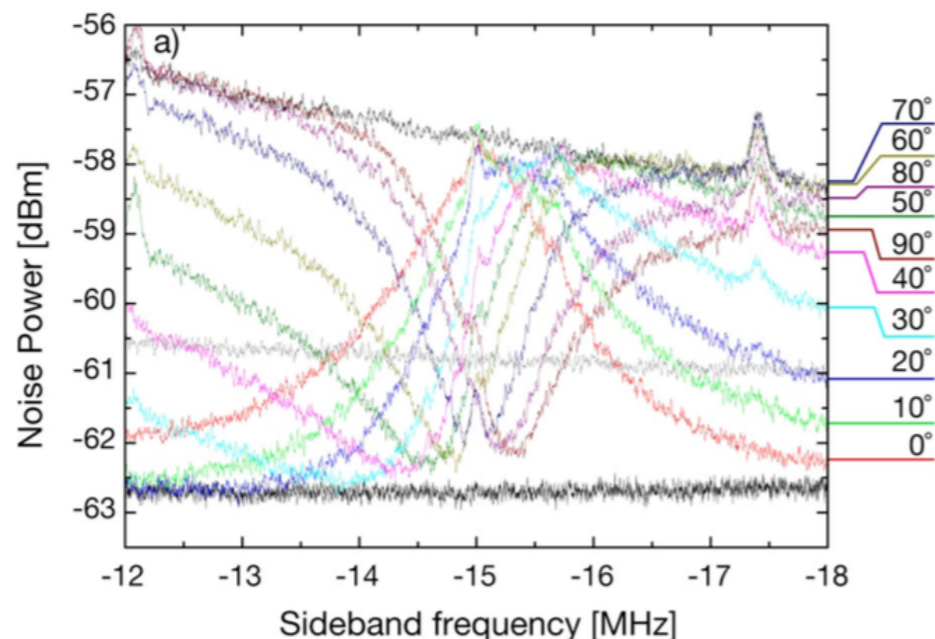
- Reflect frequency independent squeezing off a detuned Fabry-Perot cavity
- Rotation frequency depends on the cavity line-width



Optimal rotation frequency
between 40 and 70 Hz

Squeezing angle rotation already realized

@ MHz frequency

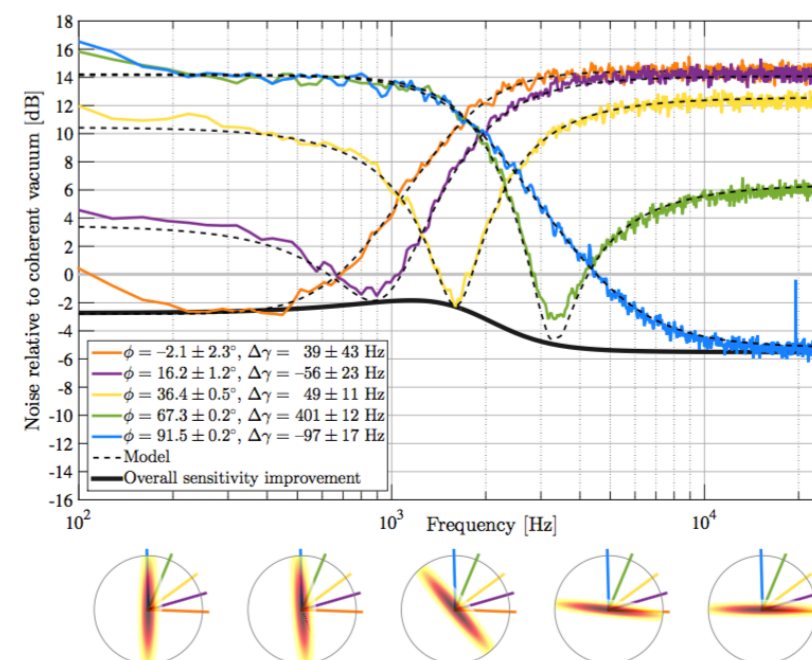


PHYSICAL REVIEW A **71**, 013806 (2005)

Experimental characterization of frequency-dependent squeezed light

Simon Chelkowski, Henning Vahlbruch, Boris Hage, Alexander Franzen, Nico Lastzka, Karsten Danzmann, and Roman Schnabel

@ kHz frequency



PRL **116**, 041102 (2016)

PHYSICAL REVIEW LETTERS

week ending
29 JANUARY 2016

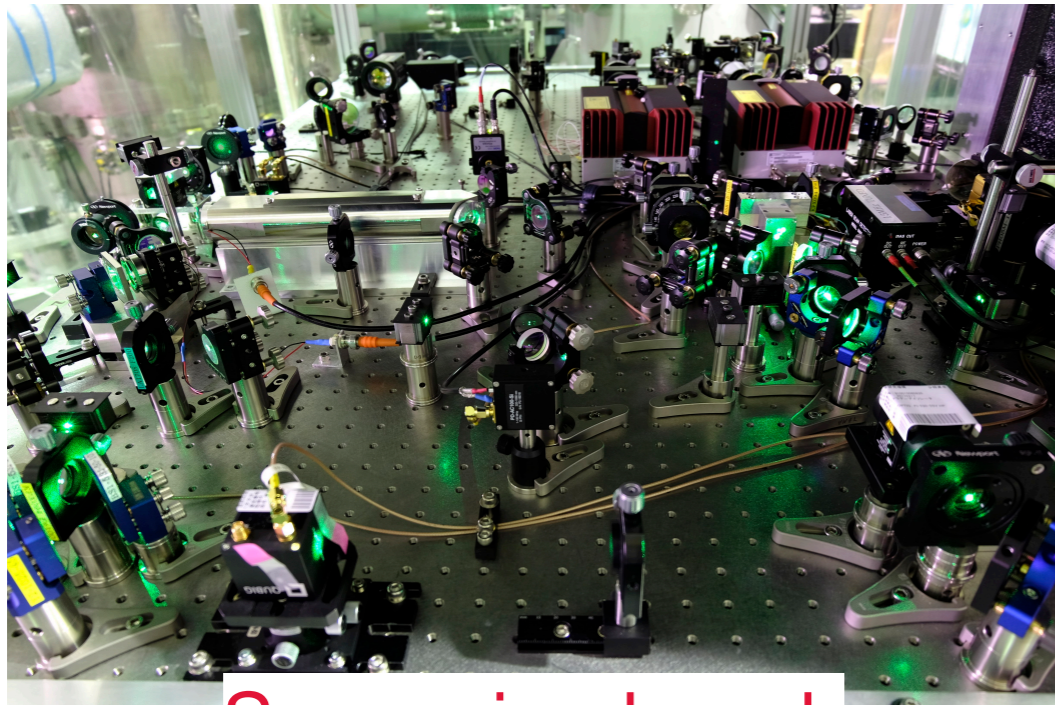
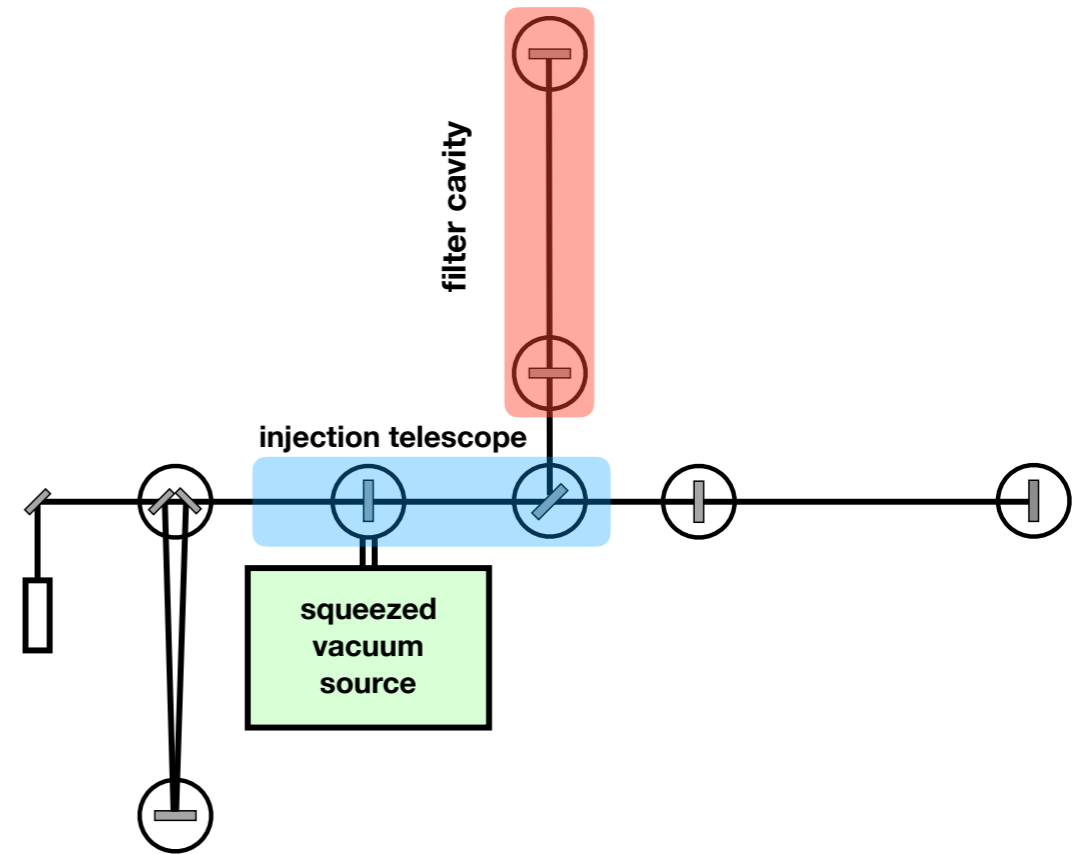
Audio-Band Frequency-Dependent Squeezing for Gravitational-Wave Detectors

Eric Oelker, Tomoki Isogai, John Miller, Maggie Tse, Lisa Barsotti, Nergis Mavalvala, and Matthew Evans*
Massachusetts Institute of Technology, Cambridge, Massachusetts 02139, USA
(Received 20 August 2015; revised manuscript received 10 December 2015; published 29 January 2016)

Our goal: **full scale filter cavity** prototype to demonstrate frequency dependent squeezing with **rotation at 70 Hz**

Experiment overview

- Cavity length: 300 m
- Finesse: 4400
- 9 dB freq. independent squeezing

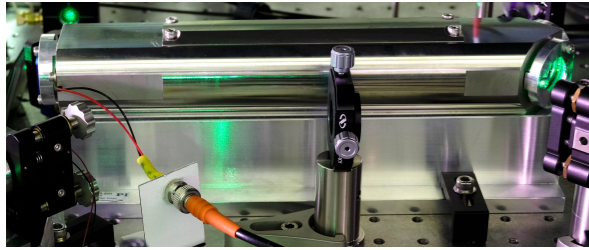


Squeezing bench

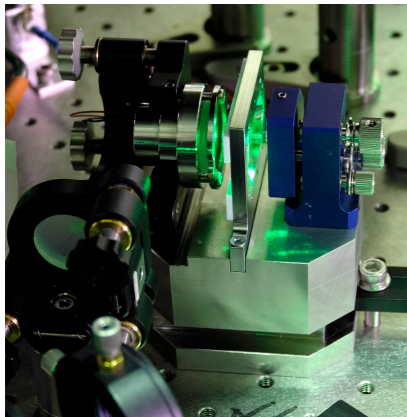


TAMA central building

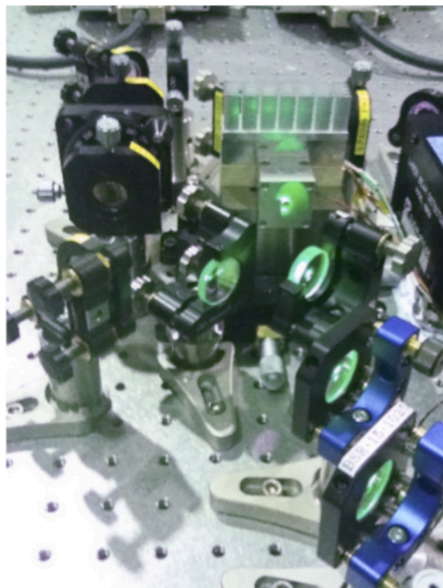
Squeezed vacuum source



Green mode cleaner and Mach-Zehnder locked



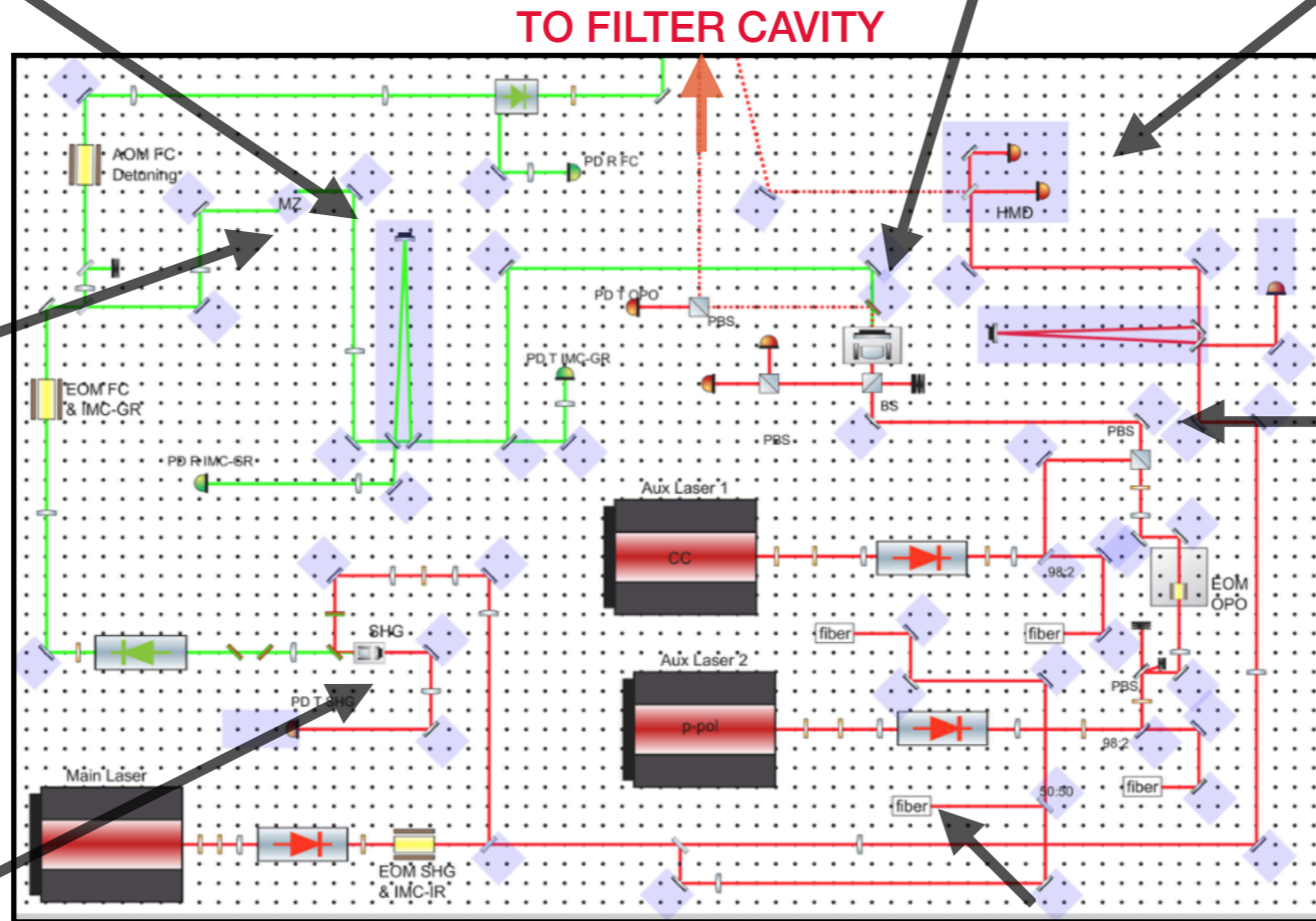
SHG assembled and operated



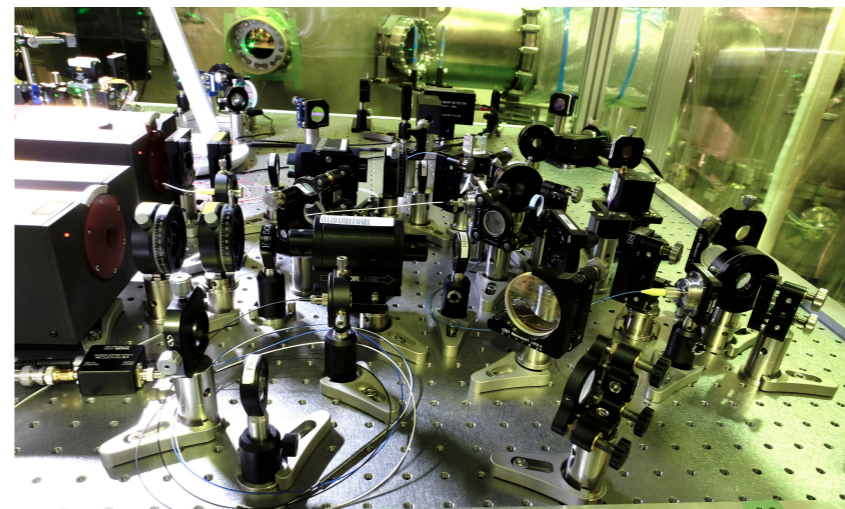
PLL for AUX lasers installed and locked

OPO assembled, alignment ongoing

Homodyne detector in progress (in collaboration with AEI, Hannover)

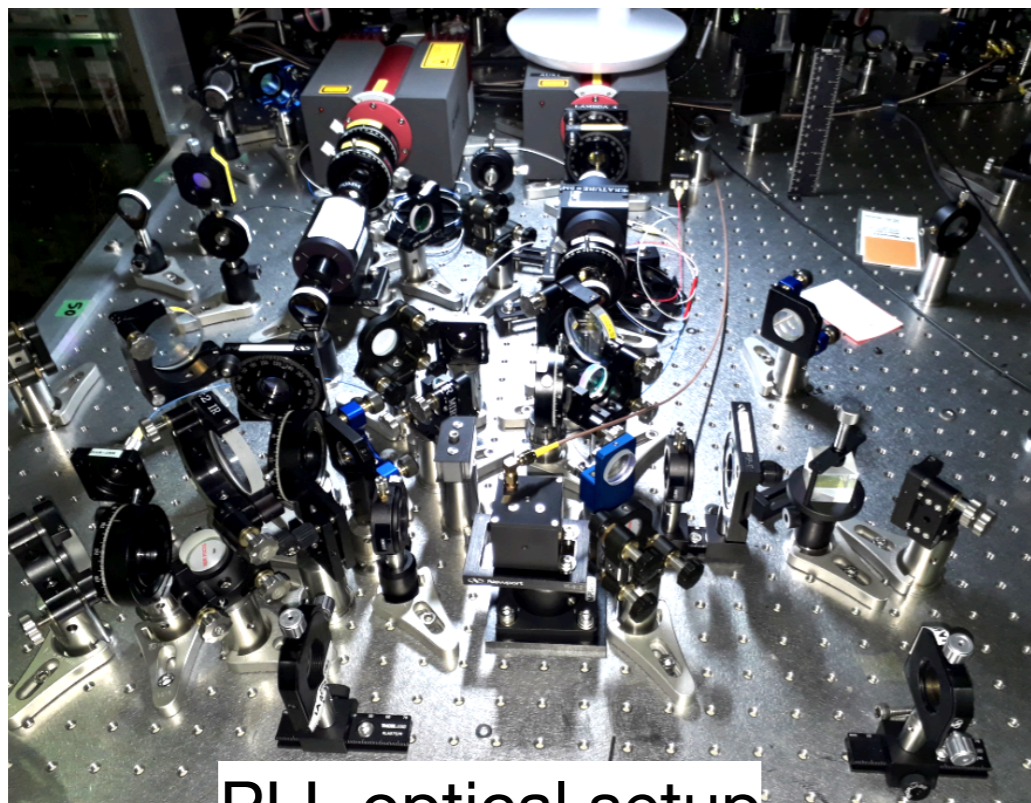


IR mode cleaner assembled

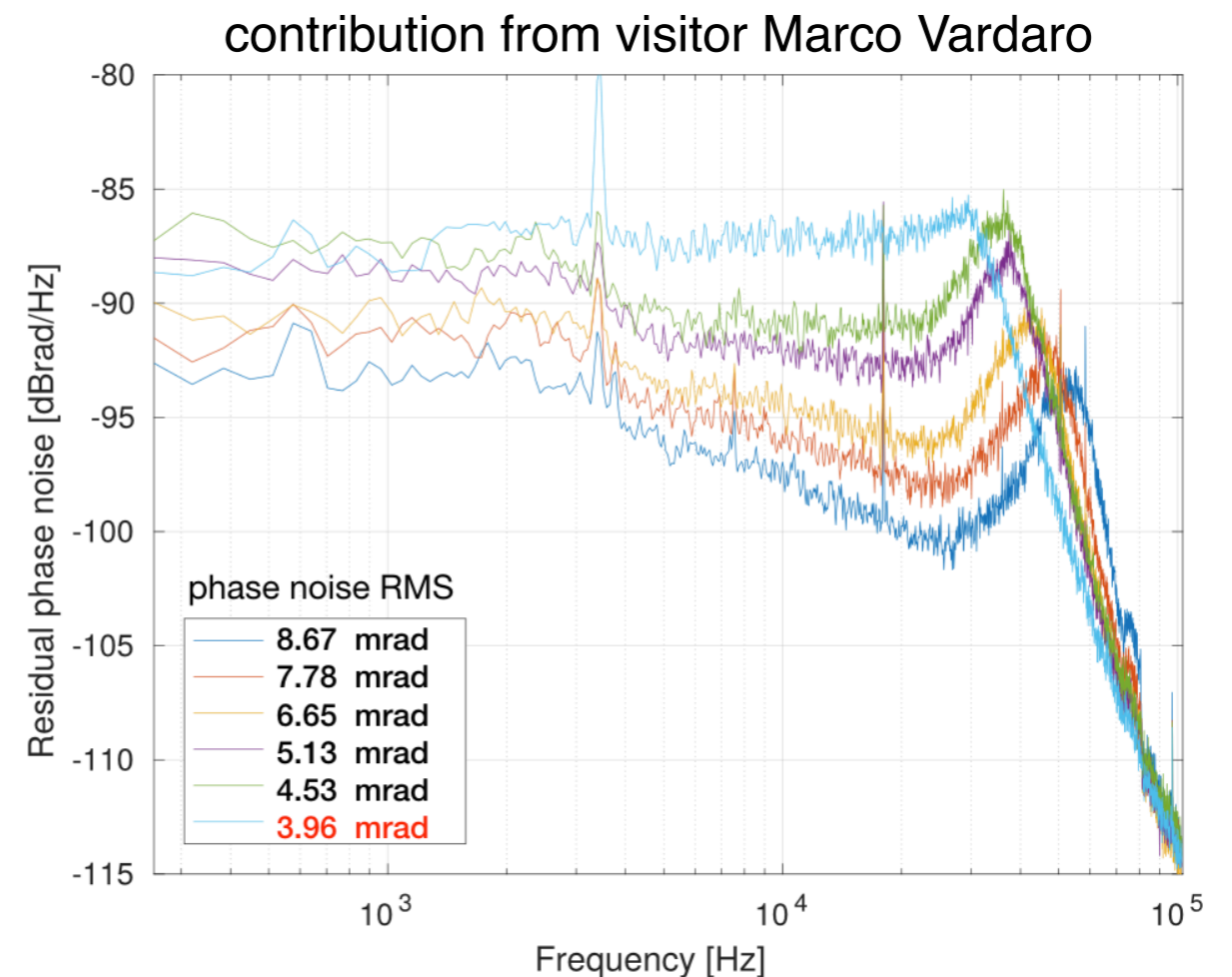


Phase lock loop (PLL) for AUX lasers

- Two AUX lasers are used for OPO length lock and control of the squeezing field phase (coherent control). They need to be phase locked to the main laser
- PLLs realized with fibered beam splitter and fiber coupled photodiode
- Electronics based on commercial Phase Frequency Detector (ADF4002)
- Residual phase noise ~ 4 mrad RMS between 100 Hz and 100 kHz

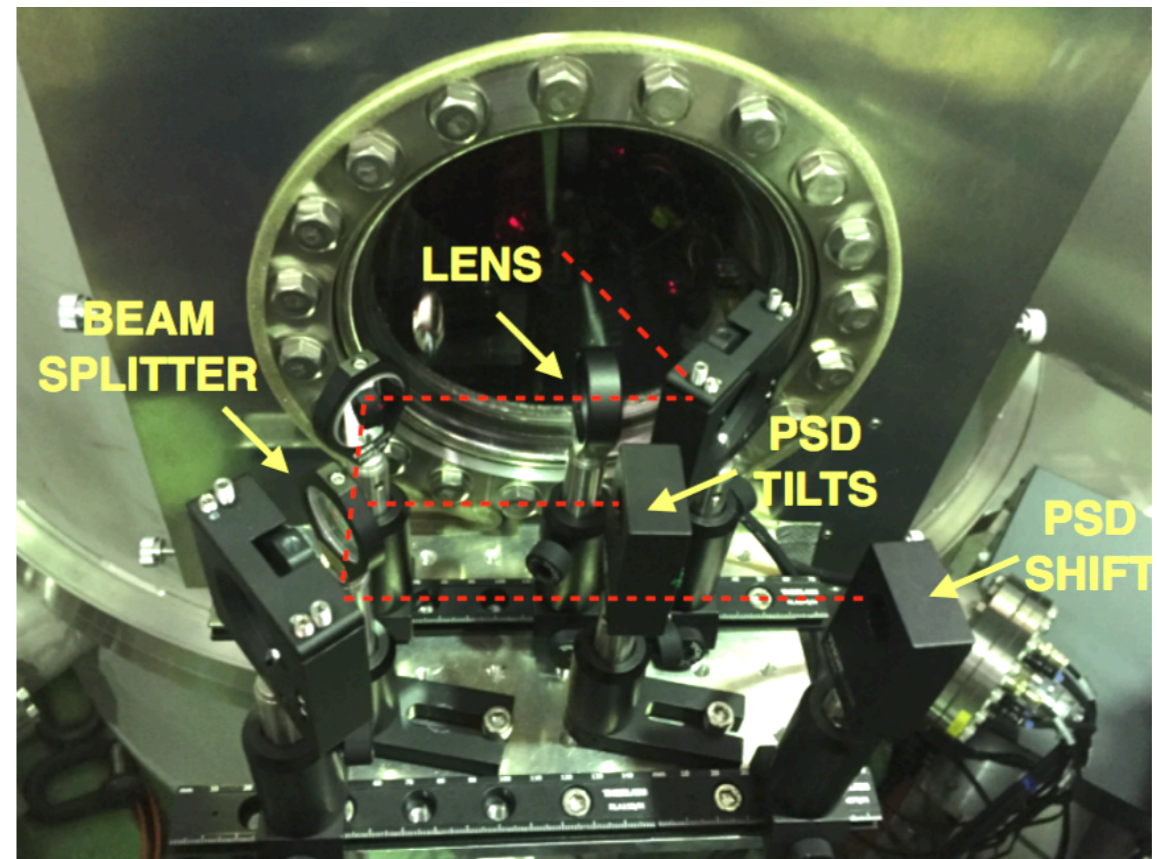
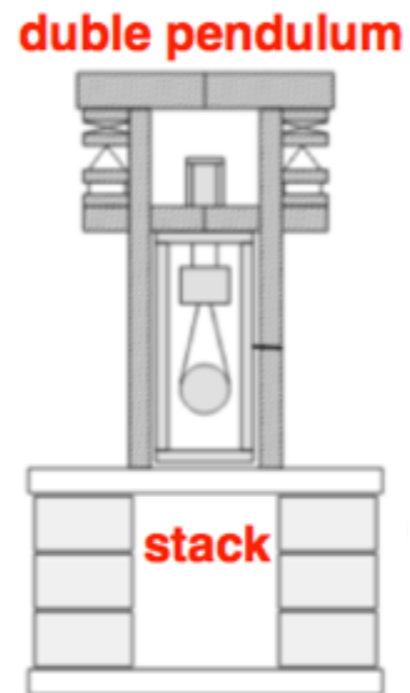
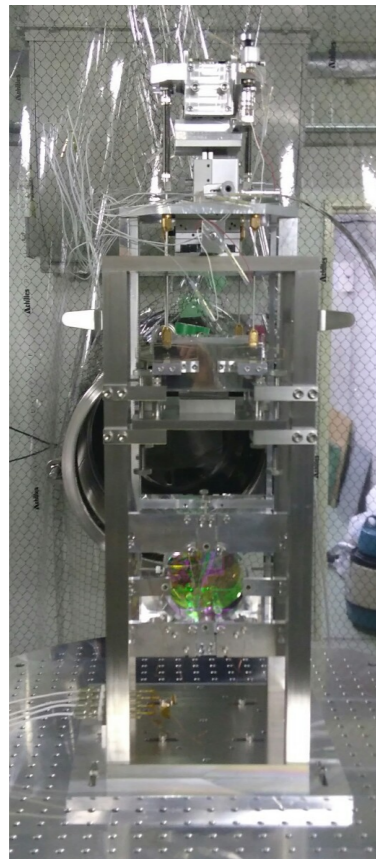


PLL optical setup



Cavity mirror suspensions

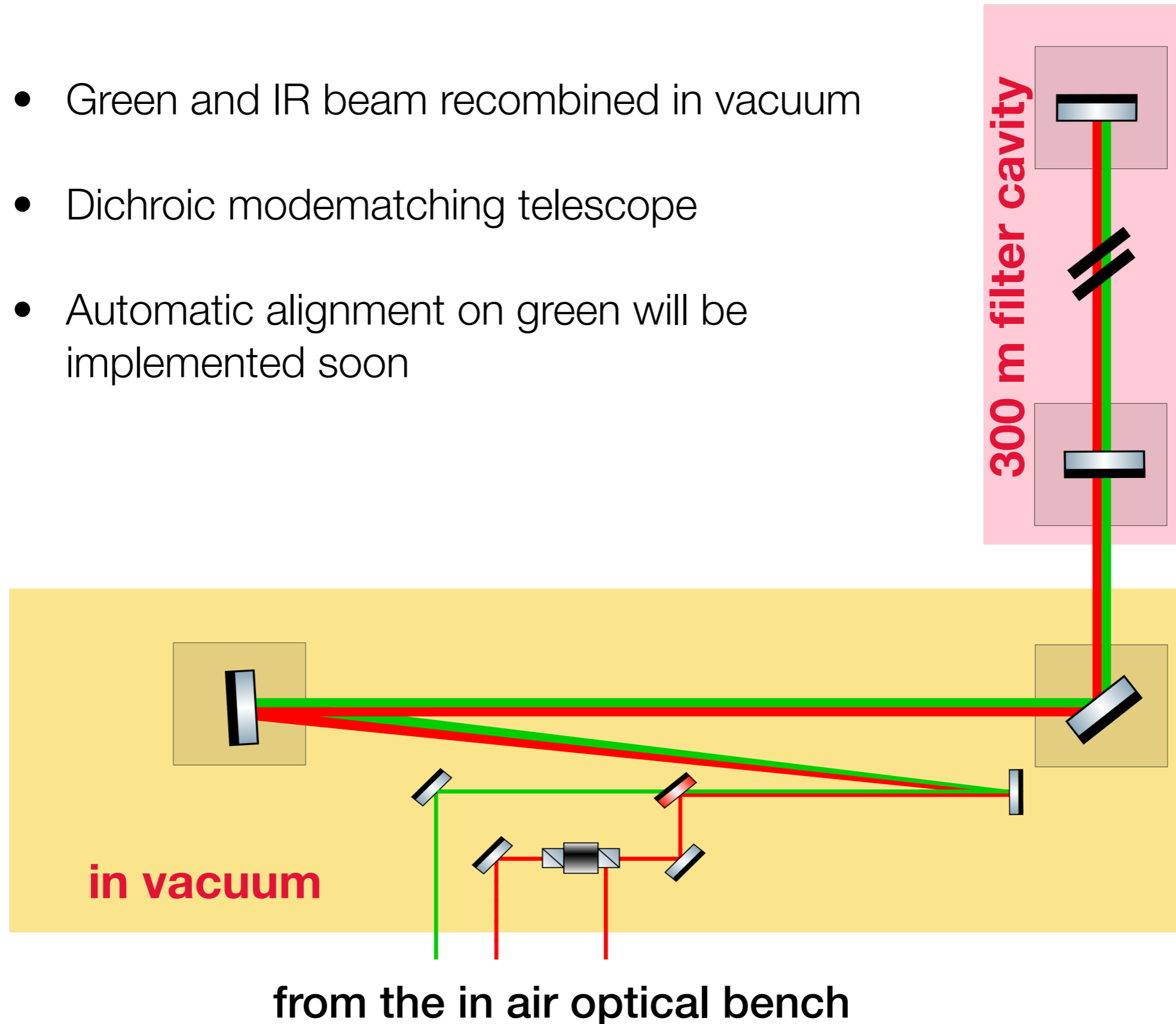
- Tama vibration isolation stack and double pendulum suspension (type C)
- Virgo-like optical levers



New control system needed to replace old TAMA LABview

Injection path and alignment

- Green and IR beam recombined in vacuum
- Dichroic modematching telescope
- Automatic alignment on green will be implemented soon

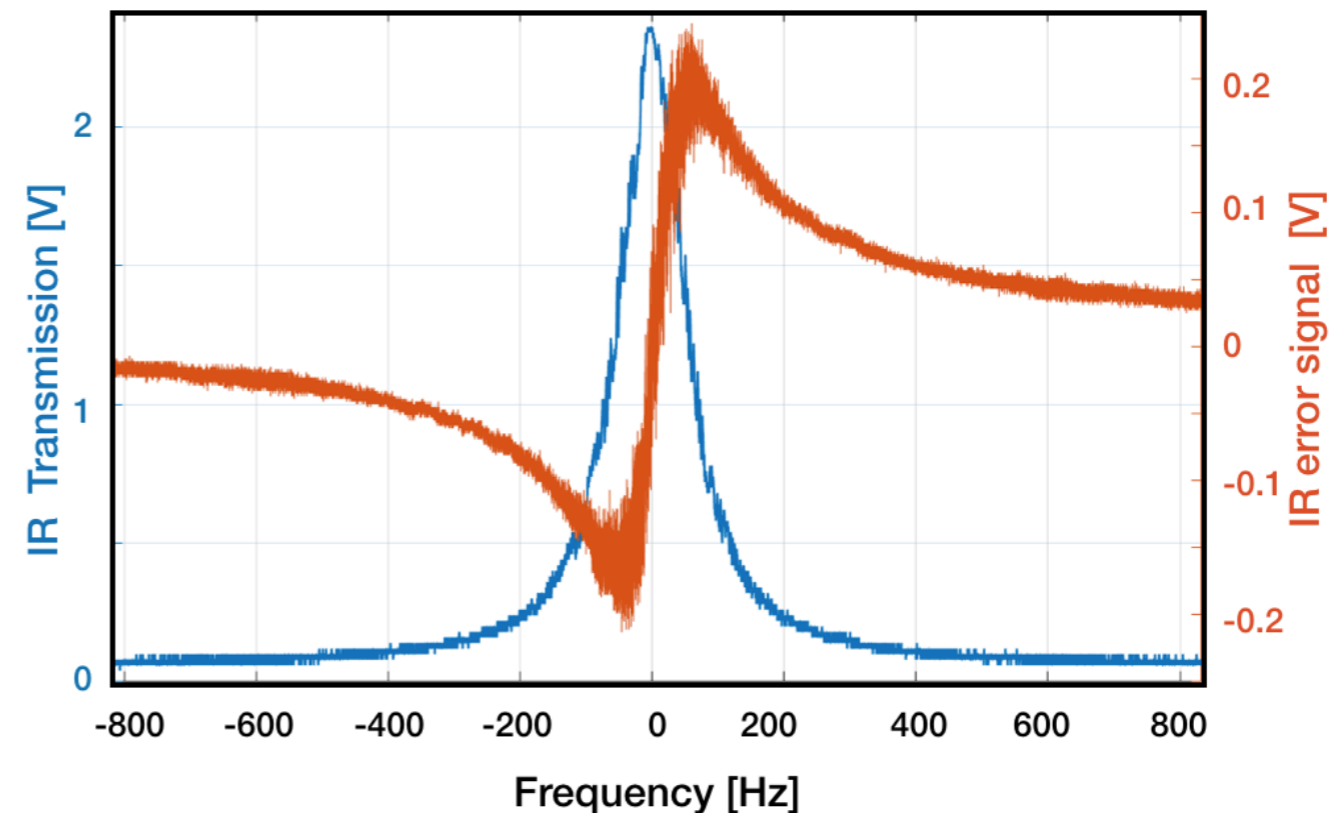


Cavity control

- Main laser locked on the cavity length using a part of the green beam from SHG.
- Analog servo with 20 kHz bandwidth (green beam finesse: 170)
- IR beam detuning controlled with a AOM on the green path

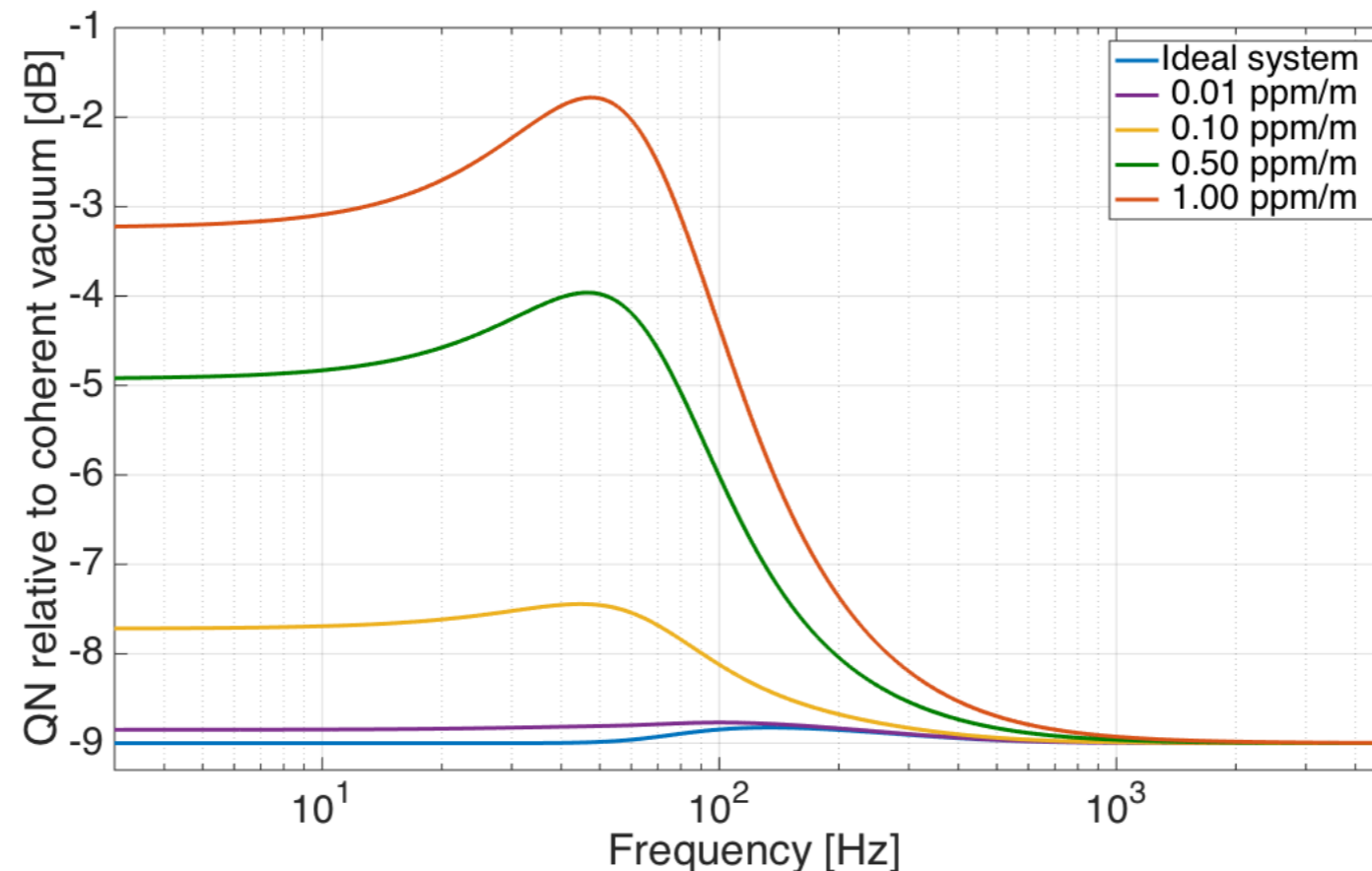


IR resonance crossing by driving AOM



Cavity characterisation: round trip losses measurement

- Important quantity for the squeezing degradation: losses per unity length
- Crucial measurement for designing filter cavities in GW detectors upgrades

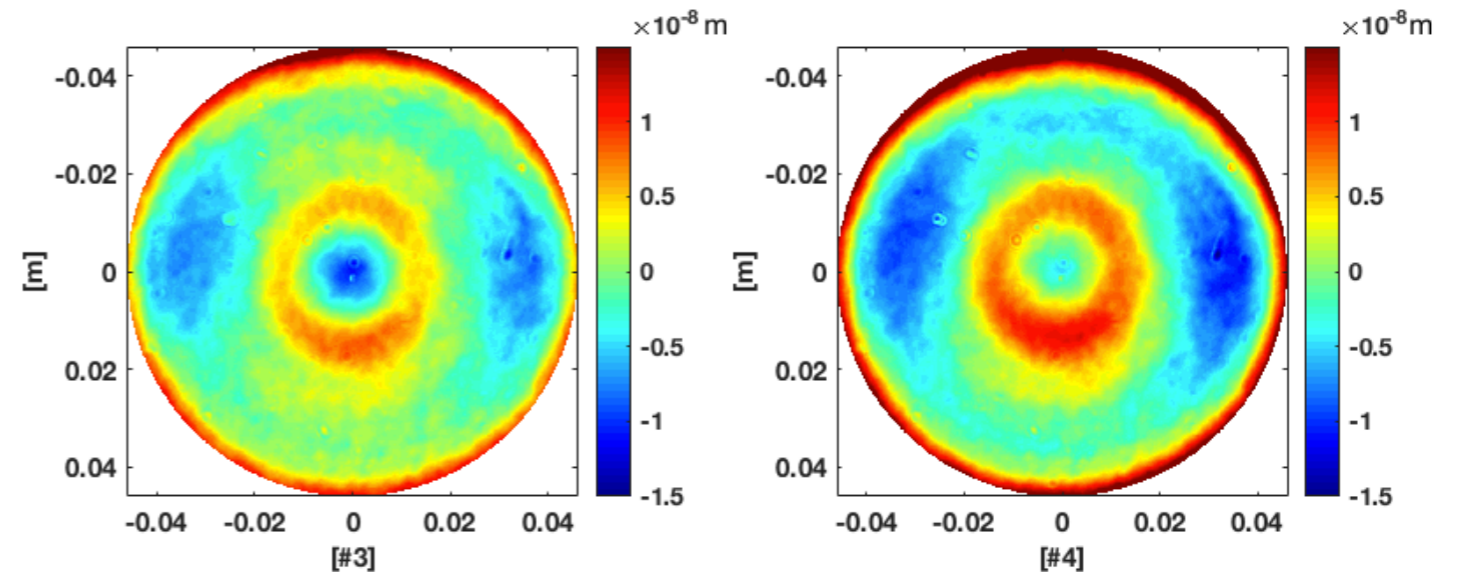
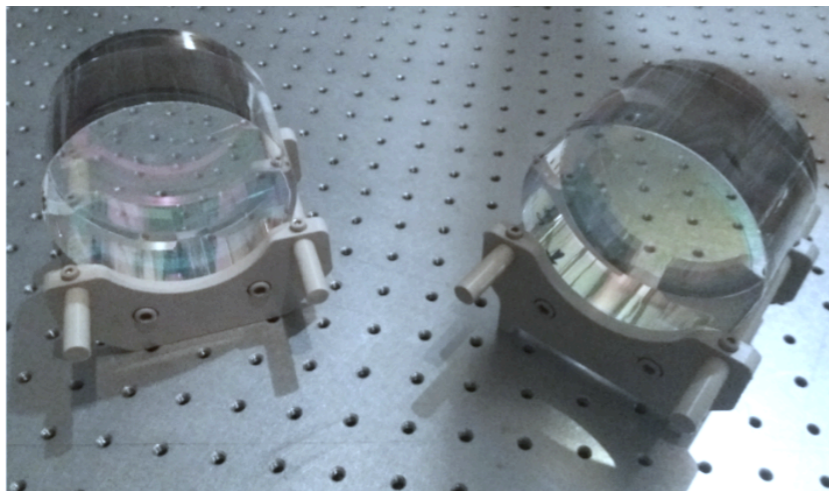
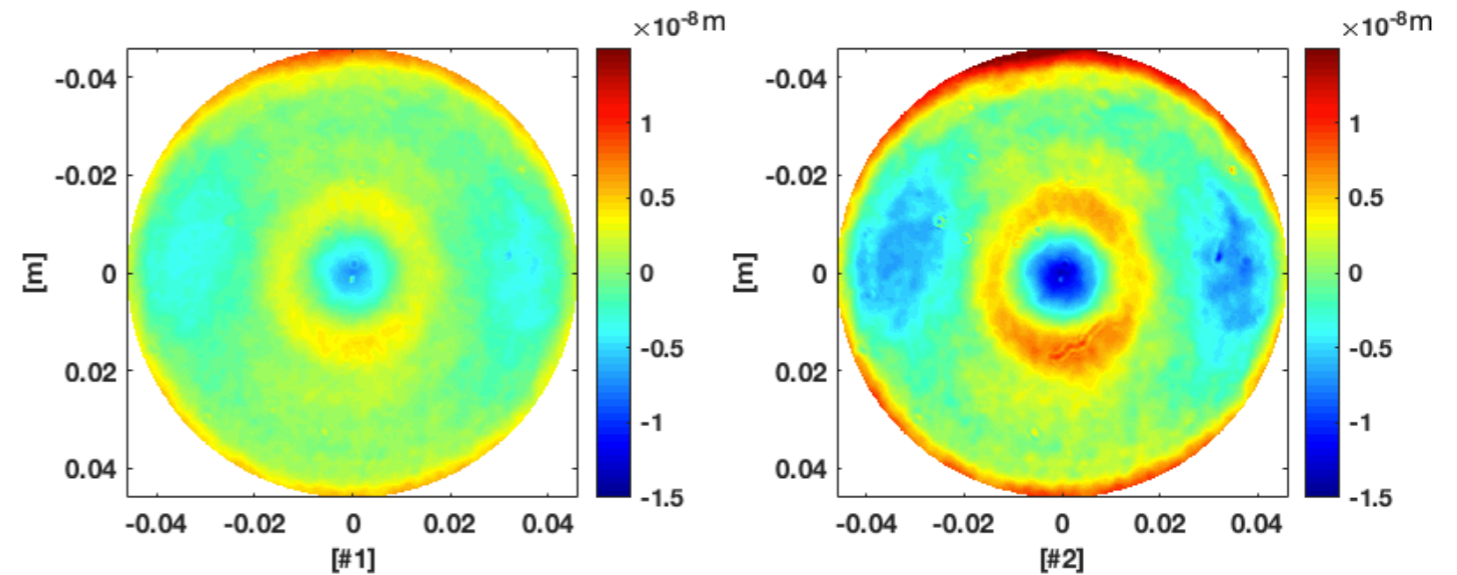


Squeezing degradation for different round trip losses values

Cavity mirrors characterisation

- Tama size: 10 cm diameter, 6 cm thickness
- Beam radius: ~ 1 cm
- Requirement on surface quality set to have 80 ppm of round trip losses

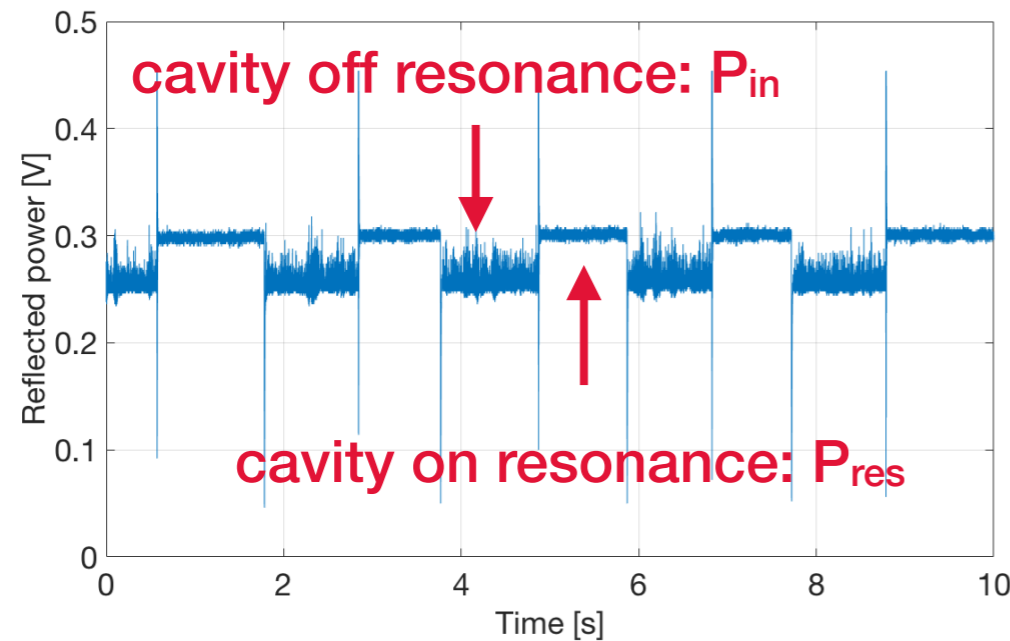
Mirror	diameter 0.05 m		diameter 0.02 m	
	RMS (nm)	PV (nm)	RMS (nm)	PV (nm)
#1	1.96	11.5	0.52	3.28
#2	2.09	12.2	0.52	3.28
#3	1.5	8.3	0.48	3.36
#4	1.94	14.8	0.48	3.28



Substrates coated and characterised at LMA

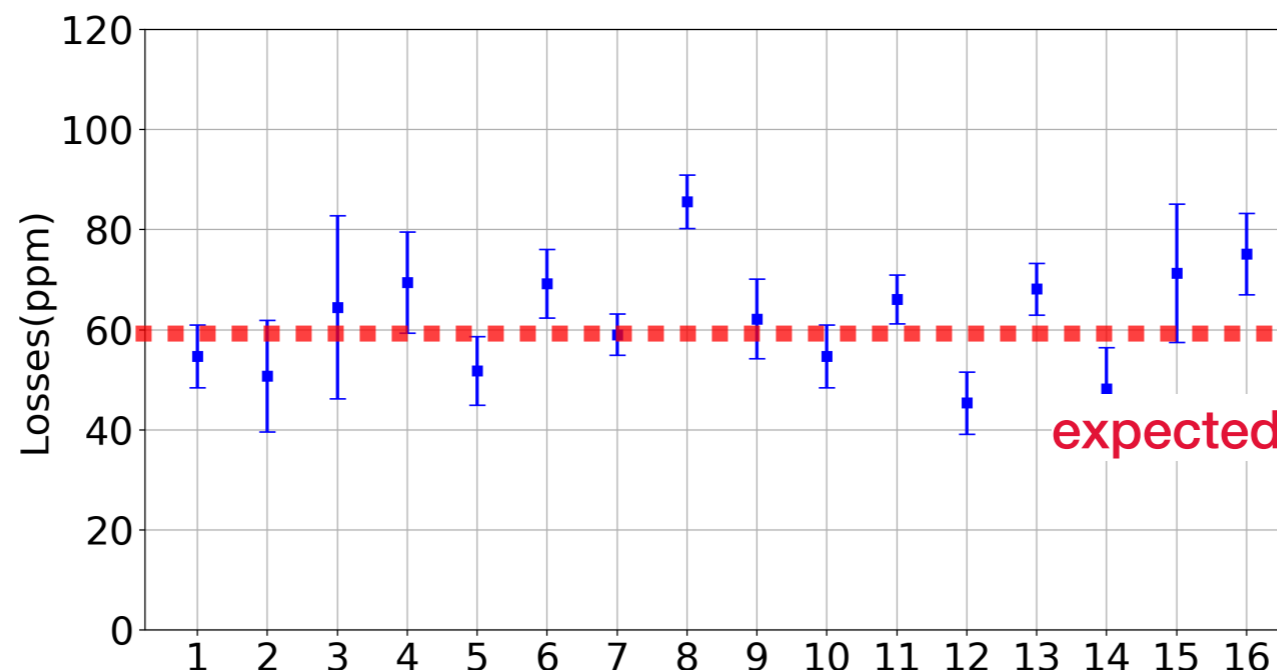
Round trip losses measurement

- Set of on/off resonances switches to measure the cavity reflectivity
- Round trip losses computed from cavity reflectivity



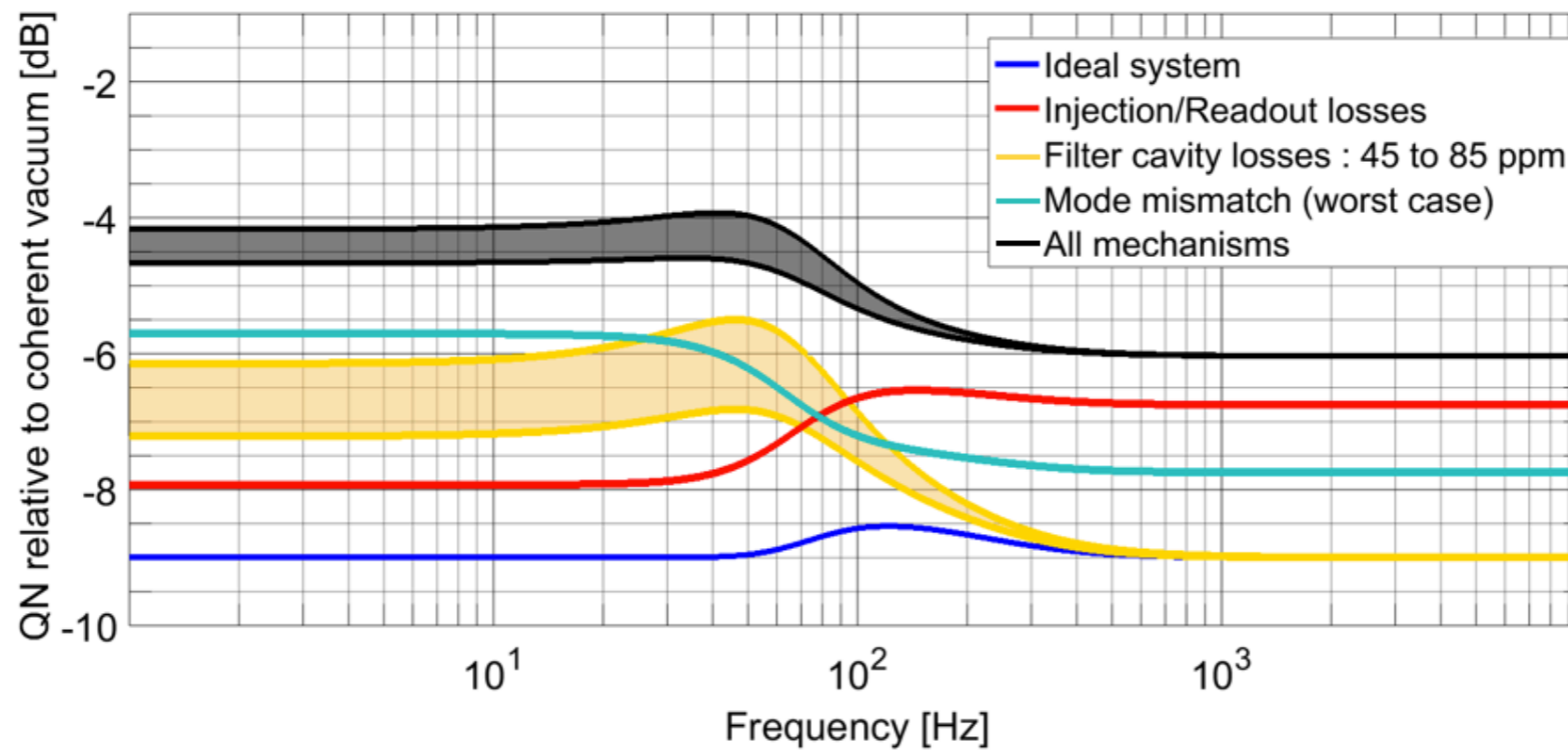
$$L \sim \frac{T_1}{2} \frac{1 - P_{\text{res}}/P_{\text{in}}}{1 + P_{\text{res}}/P_{\text{in}}}$$

- **Round trip losses between 45 and 85 ppm**



Expected squeezing level

- Round trip losses below 85 ppm are allows for ~ 4 dB of squeezing at low frequency (assuming 9 dB of injected squeezing)



PHYSICAL REVIEW D **98**, 022010 (2018)

Measurement of optical losses in a high-finesse 300 m filter cavity for broadband quantum noise reduction in gravitational-wave detectors

Eleonora Capocasa,^{1,2,*} Yuefan Guo,³ Marc Eisenmann,⁴ Yuhang Zhao,^{1,5} Akihiro Tomura,⁶ Koji Arai,⁷ Yoichi Aso,¹ Manuel Marchiò,¹ Laurent Pinard,⁸ Pierre Prat,² Kentaro Somiya,⁹ Roman Schnabel,¹⁰ Matteo Tacca,¹¹ Ryutaro Takahashi,¹ Daisuke Tatsumi,¹ Matteo Leonardi,¹ Matteo Barsuglia,² and Raffaele Flaminio^{4,1}

(Received 27 May 2018; published 31 July 2018)

Visitors and new collaboration

- Many contributions given from visitors



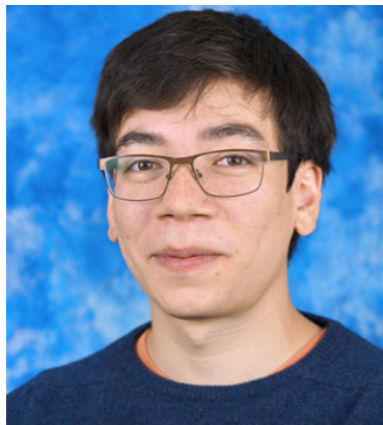
Emil Schreiber
GEO600/AEI



Marco Vardaro
Padova University



Matteo Barsuglia
APC/CNRS



Marc Eisenmann
LAPP/CNRS



Yuefan Guo
Beijing Normal University



Pierre Prat
APC/CNRS

Visitors and new collaboration

- Visit of Ray-Kuang Lee's quantum optics group from National Tsing Hua University, Taiwan
- Discussion about future collaboration with them



國立清華大學
NATIONAL TSING HUA UNIVERSITY

Summary and perspectives

Status

- Filter cavity installed and controlled
- Cavity round trip losses between 45 and 85 ppm
- Squeezed vacuum source integration almost completed

Next steps

- Operation of the squeezing vacuum source before the end of FY2018
- Cavity automatic alignment implementation
- Upgrade of the digital control system
- Final step: injection of freq. independent squeezing into the filter cavity

More information can be found on our wiki page:

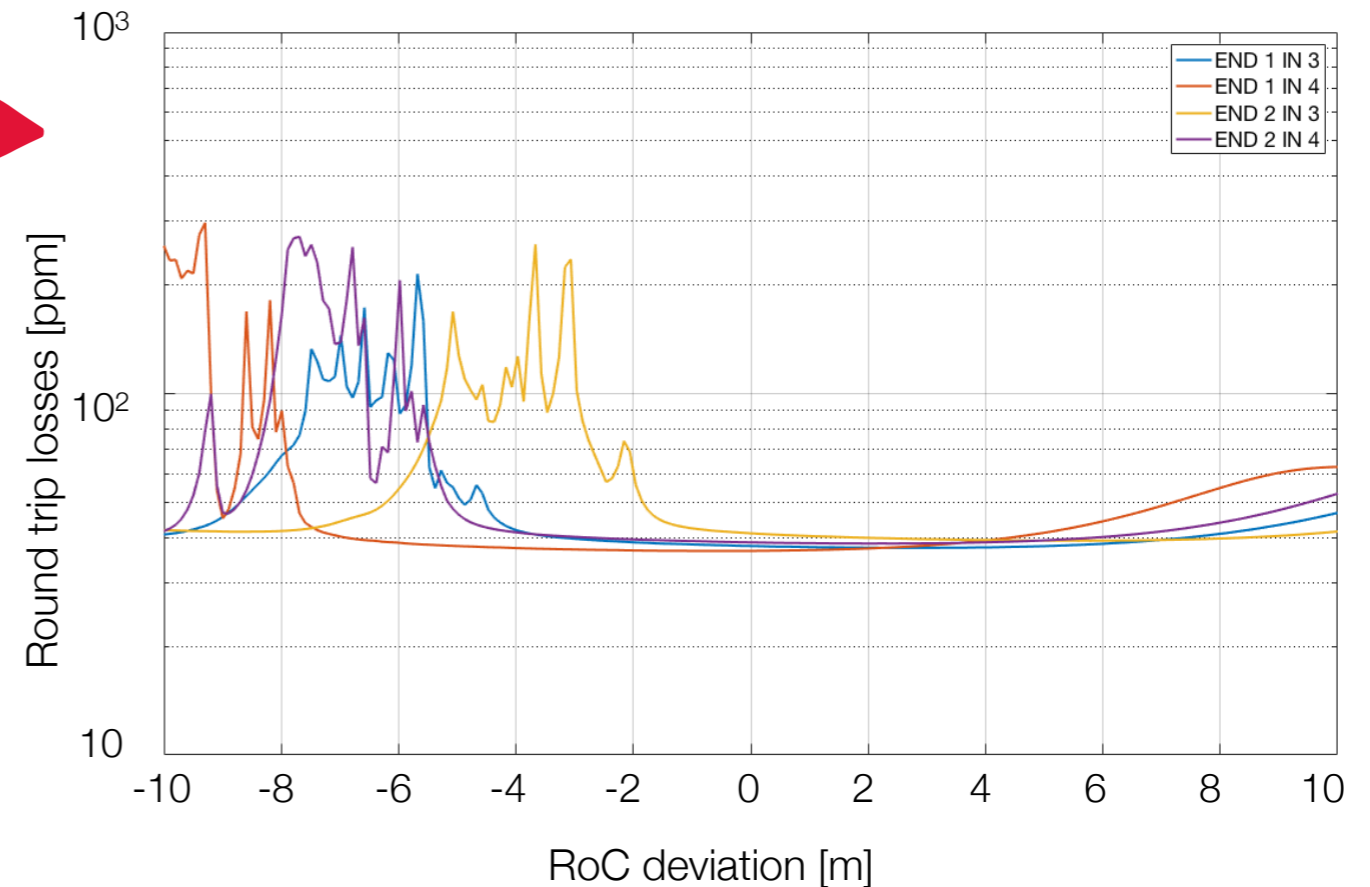
<https://gwpo.nao.ac.jp/wiki/FilterCavity>

EXTRA SLIDES

Round trip losses budget

RTL REQUIREMENT : 80 ppm

- ~ **40 ppm** from flatness (simulation)
- ~ **15 ppm** from roughness and point defect (measured)
- ~ **5 ppm** from absorption and transmission (measured)



TOTAL EXPECTED RTL : ~ 60 ppm

How to measure round trip losses?

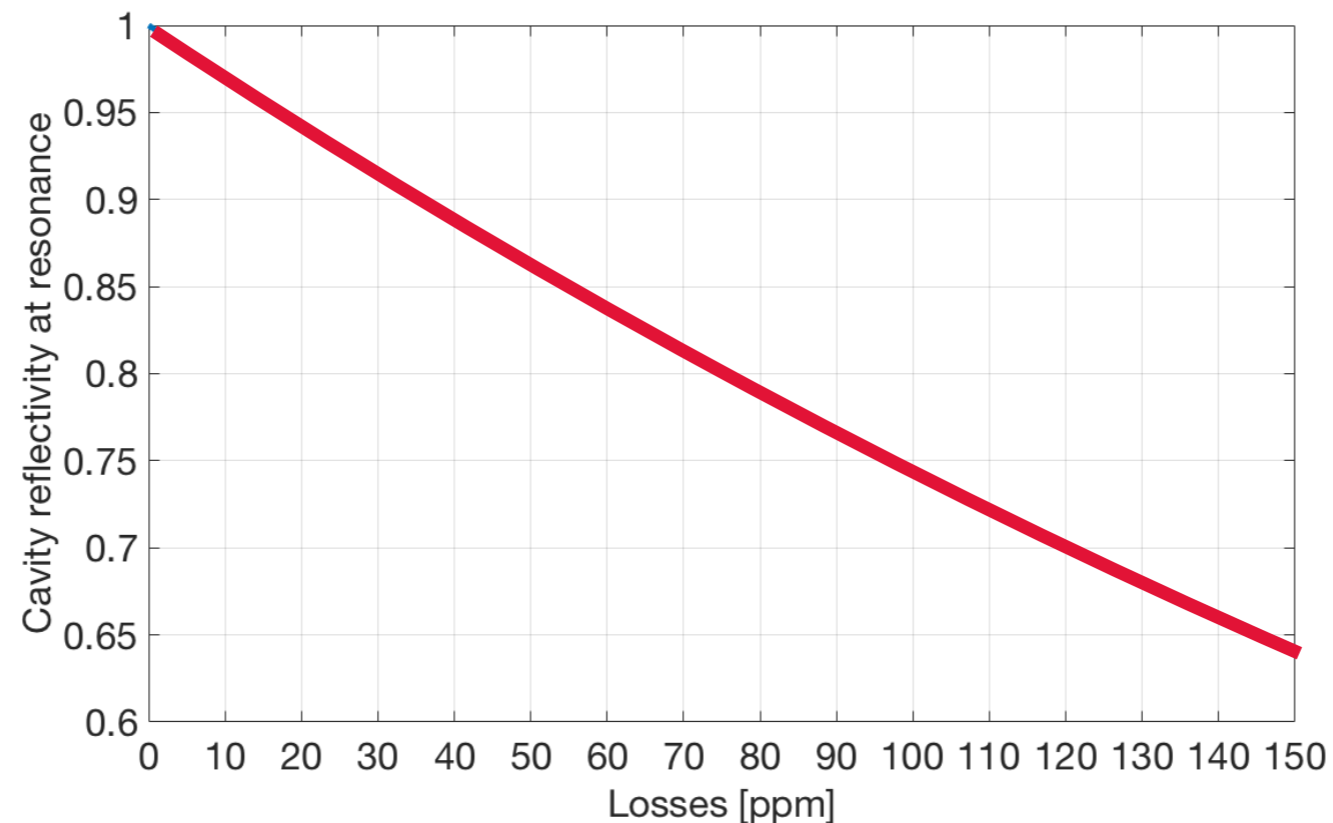
- From the cavity reflectivity at resonance
- Reflectivity is less affected from the input mirror transmissivity (with respect to finesse or decay time)

$$R_{\text{cav}} = \left[\frac{r_1 - r_2}{1 - r_1 r_2} \right]^2$$

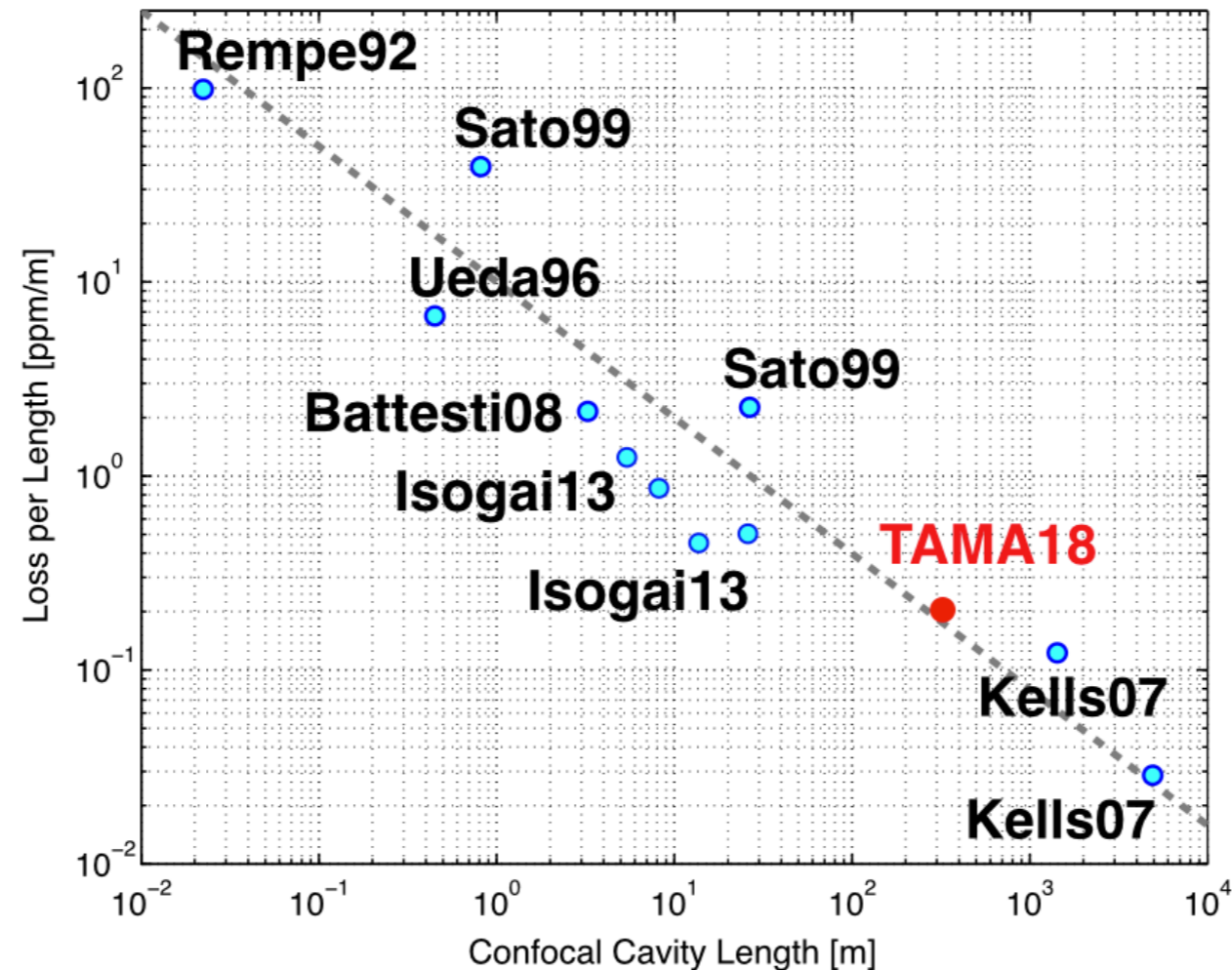


$$r_2 = \sqrt{1 - T_2 - L} \simeq \sqrt{1 - L}$$

$$L = \frac{T_1}{2} \cdot \frac{1 - R_{\text{cav}}}{1 + R_{\text{cav}}}$$



Comparison with round trip losses in literature



$$L_{\text{rt}}(\mathcal{L}_{\text{confocal}}) = 10 \text{ ppm} \cdot \left(\frac{\mathcal{L}_{\text{confocal}}}{1 \text{ m}} \right)^{0.3}$$

$\mathcal{L}_{\text{confocal}}$ is the length of the confocal cavity which has the same spot size at its mirrors as the cavity whose losses are reported

PHYSICAL REVIEW D **88**, 022002 (2013)

Realistic filter cavities for advanced gravitational wave detectors

M. Evans, L. Barsotti, and P. Kwee

Massachusetts Institute of Technology, Cambridge, Massachusetts 02139, USA

J. Harms

INFN, Sezione di Firenze, Sesto Fiorentino 50019, Italy

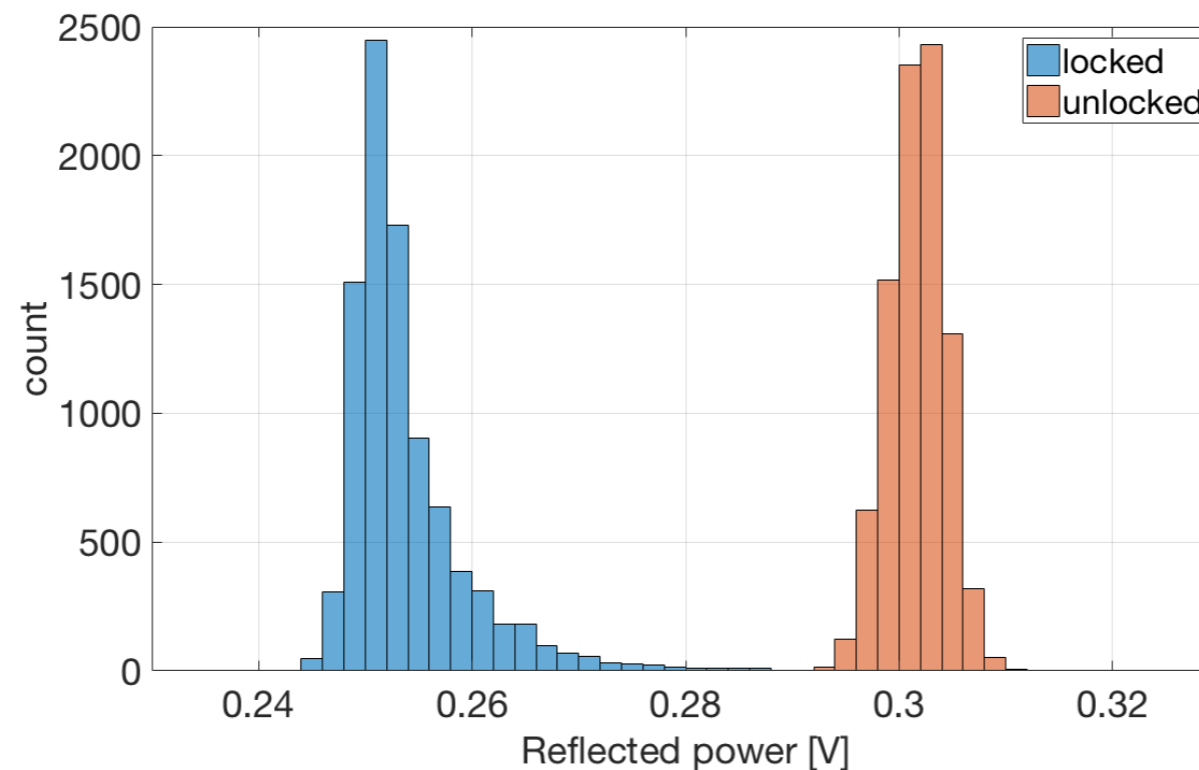
H. Miao

California Institute of Technology, Pasadena, California 91125, USA

(Received 9 May 2013; published 29 July 2013)

How to estimate the reflected power

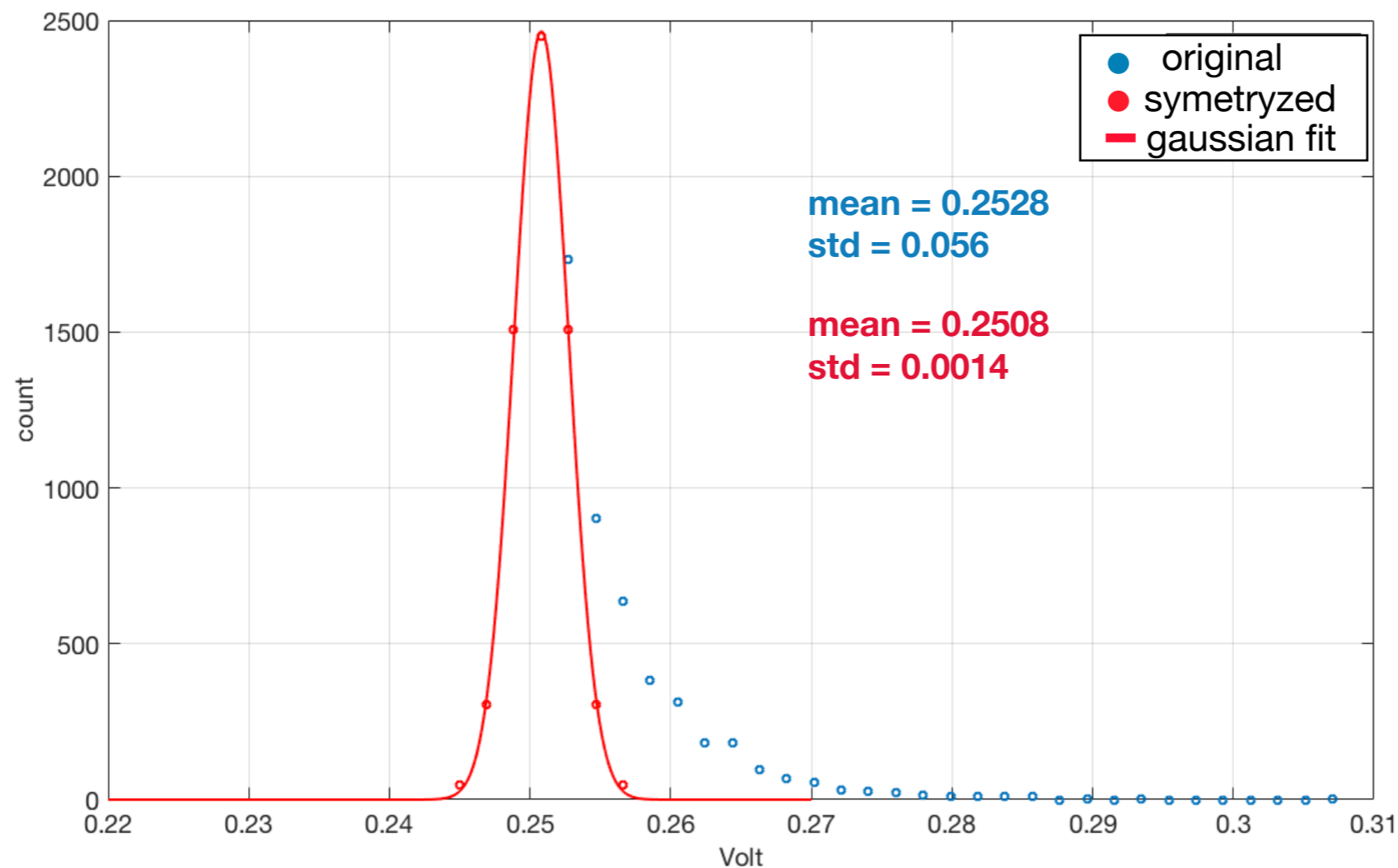
- Power reflected has some fluctuation which show different features if the cavity is locked or unlocked



- Cavity unlocked: Gaussian histogram. Main influence: input power fluctuations
- Cavity unlocked: asymmetric distribution. Influence of the input power fluctuations, cavity alignment fluctuations, finite lock accuracy

How to estimate the reflected power

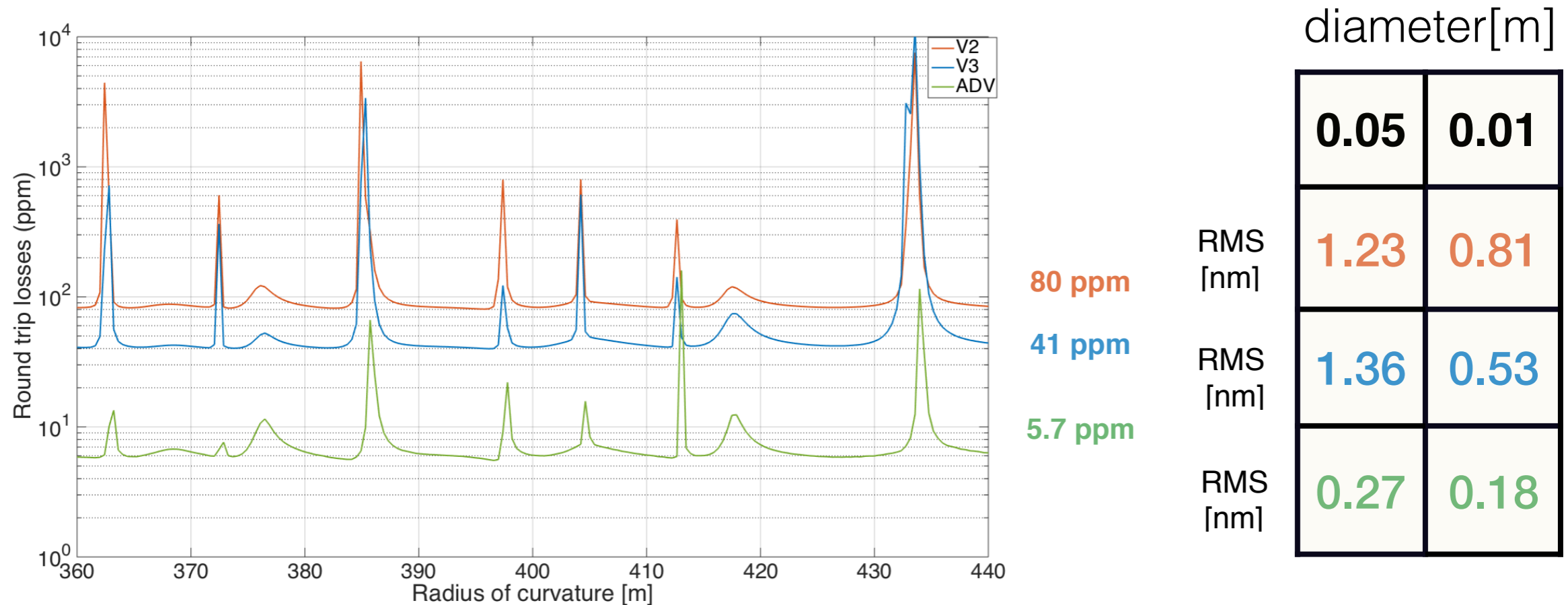
- Since the fluctuation which tends to increase the power (e.g. those of the lock accuracy) can bring to an underestimation of the losses) **we decided to consider the mode of the fluctuation histograms instead of the mean**



- The difference in the estimated losses is < 5 ppm. Results is not strongly dependent on this choice

Cavity design: requirement on the mirrors flatness

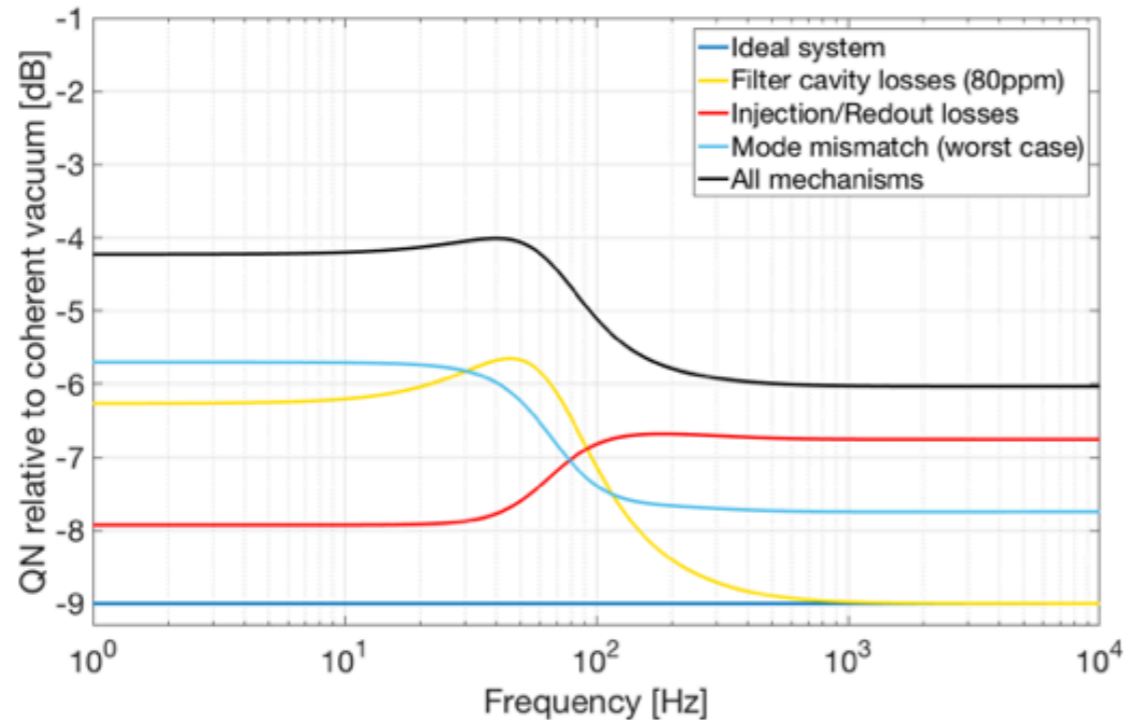
FFT simulation to measure RTL using real Virgo Mirror map



In order to set the RTL threshold we need to consider the other squeezing degradation sources

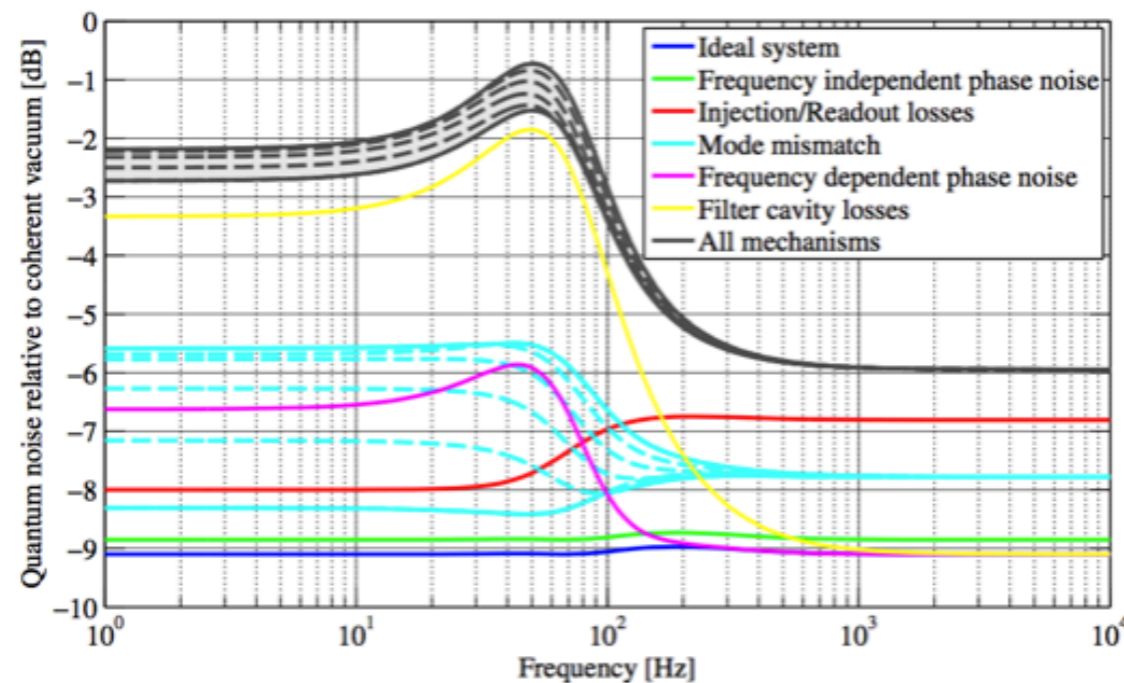
Squeezing degradation budget

TAMA 300 filter cavity



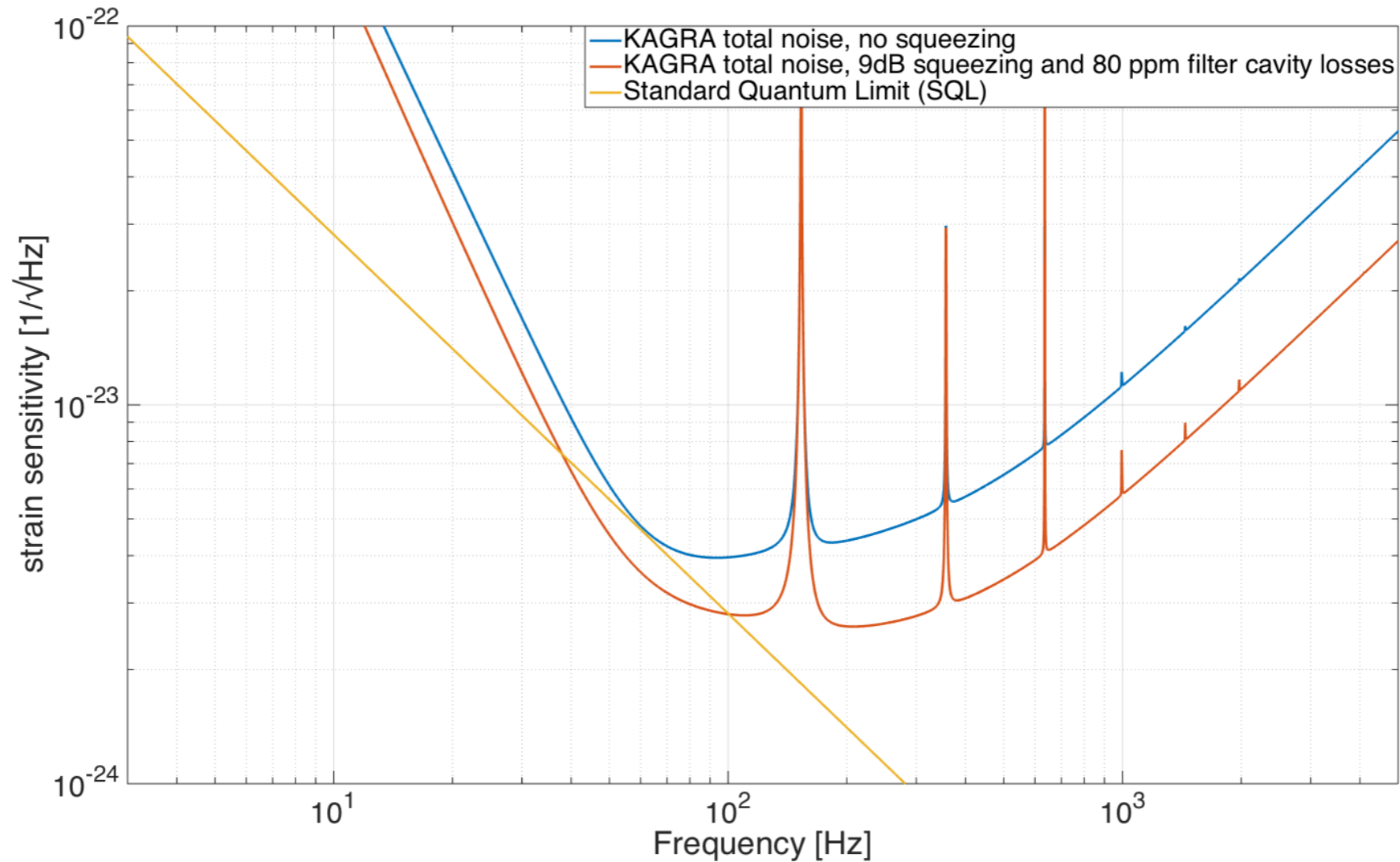
length	16 m 300 m
RTL00	16ppm 40ppm
injection losses	5%
readout losses	5%
squeezer-filter cavity mismatch	2%
squeezer-local oscillator mismatch	5%
δL (rms)	0.3 pm

16 m filter cavity at MIT



- **Important quantity: ppm per meter**
- RTL of 80 ppm (corresponding to Virgo mirrors quality) are low enough
- RTL of 6 ppm (corresponding to AdVirgo mirrors quality) makes degradation from cavity losses completely negligible

Expected improvement on KAGRA sensitivity



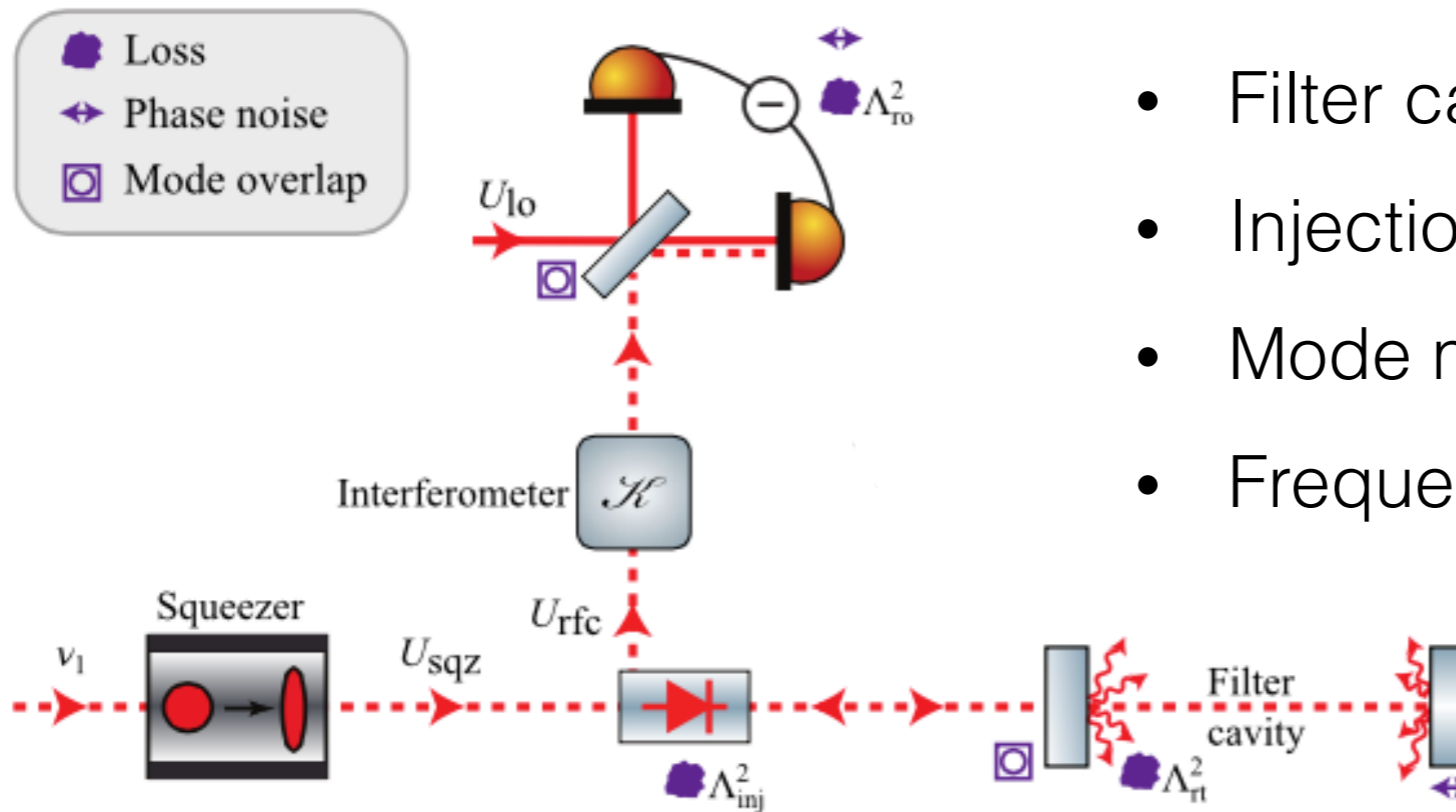
Horizon without squeezing

BNS = 360 Mpc BBH = 3.28 Gpc

Horizon with squeezing

BNS = 509 Mpc BBH = 4.42 Gpc

Squeezing degradation sources



- Filter cavity losses
- Injection/readout losses
- Mode mismatch
- Frequency-dependent phase noise

P.Kwee et al. "Decoherence and degradation of squeezed states in quantum filter cavities" Phys. Rev. D 90 062006 (2014)

Quantum fluctuation entering with the losses should be taken into account

$$N(\zeta) = |\bar{\mathbf{b}}_{\zeta} \cdot \mathbf{T}_1 \cdot v_1|^2 + |\bar{\mathbf{b}}_{\zeta} \cdot \mathbf{T}_2 \cdot v_2|^2 + |\bar{\mathbf{b}}_{\zeta} \cdot \mathbf{T}_3 \cdot v_3|^2$$

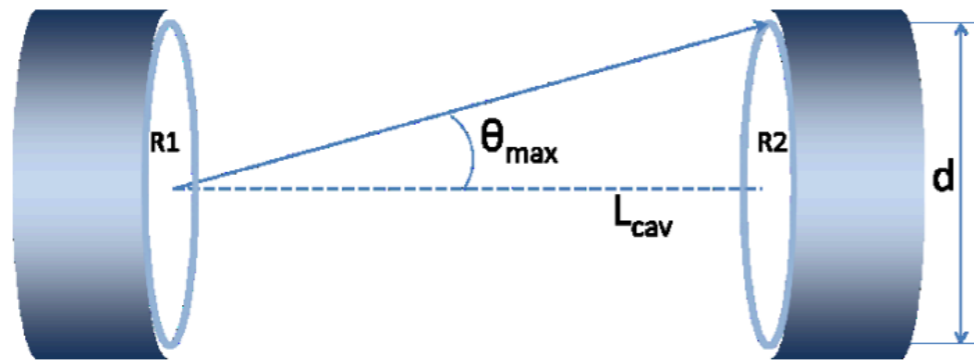
↑
squeezed
field

↑
vacuum fluctuation
due to losses
before the ITF

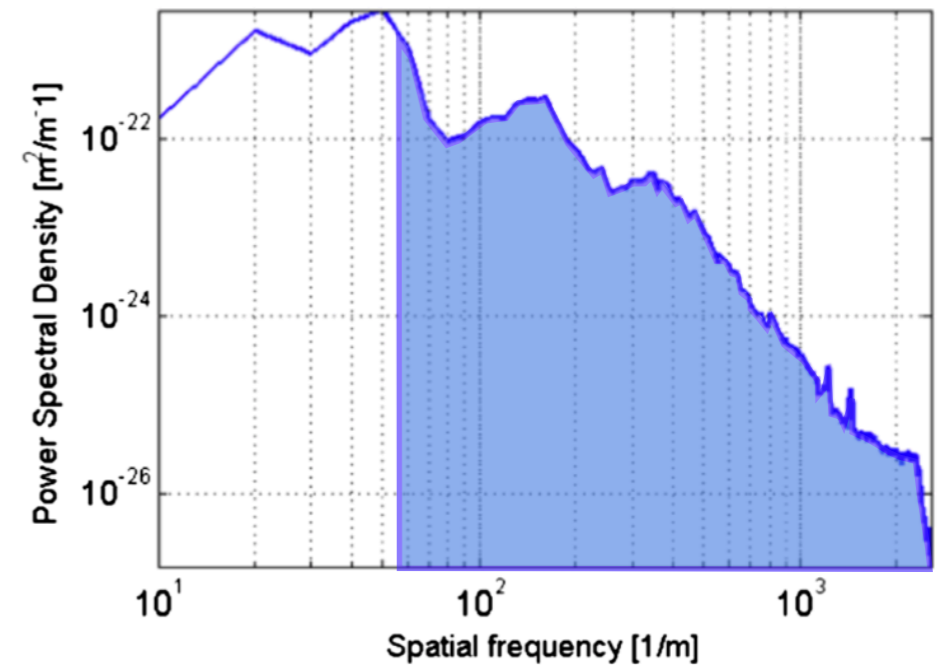
↑
vacuum fluctuation
due to losses after
the ITF

Scattering from mirror defects

Diffraction angle: $\theta = \lambda \times f$



$$f_{\text{limit}} = \frac{d}{2L \times \lambda}$$



$$\sigma^2 = \int_{f_1}^{f_2} PSD(f) df$$

- Flatness: $10 \text{ m}^{-1} - 10^3 \text{ m}^{-1}$
- Roughness: $10^3 \text{ m}^{-1} - 10^5 \text{ m}^{-1}$
- Point defects $> 10^5 \text{ m}^{-1}$

Scattering golden rule:

$$\text{losses}_{(f > f_{\text{limit}})} = \left(\frac{4\pi \times \sigma}{\lambda} \right)^2$$

Why ppm/meter are important?

Total losses: RTL per number of round trip $N \sim 1/T_f$

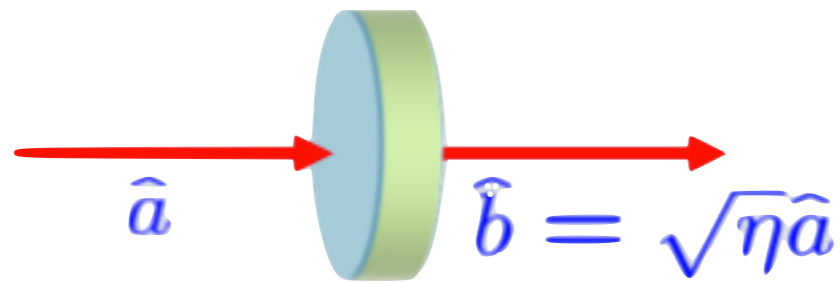
$$\mathcal{E} \approx \frac{\epsilon}{T_f}$$

Optimal rotation: filter Cavity bandwidth comparable with ITF bandwidth γ

$$T_f \approx \frac{4\gamma L_f}{c} \quad \longrightarrow \quad \mathcal{E} \approx \frac{c\epsilon}{4\gamma L_f} \propto \frac{\epsilon}{L_f}$$

Optical losses degrade squeezing

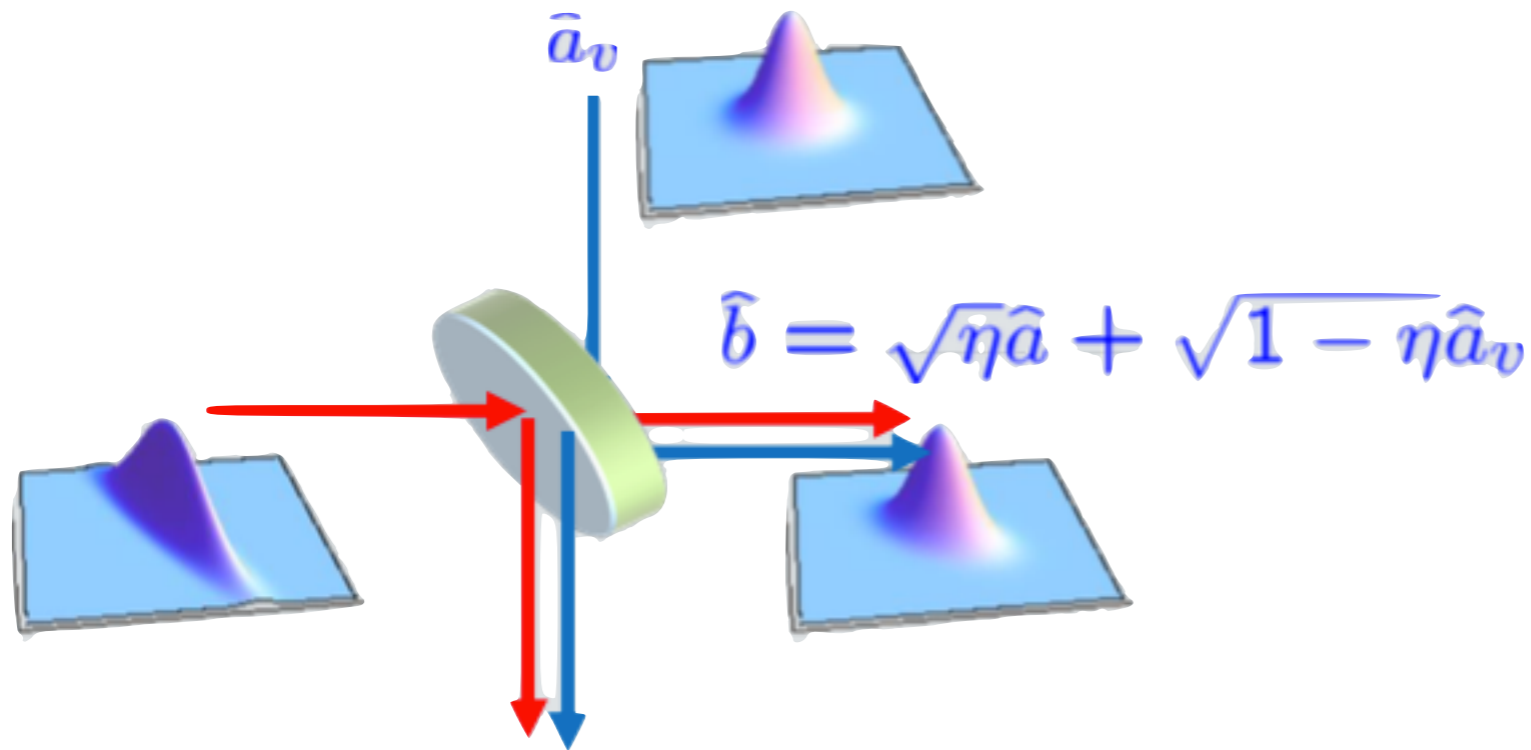
- Naive model



$$[\hat{a}, \hat{a}^\dagger] = 1$$

$$[\hat{b}, \hat{b}^\dagger] = \eta \neq 1$$

- Consistent model



Squeezing degraded
because of its
recombination with
non squeezed
vacuum